



POLYTECH.MONS

# BELVAC/NEVAC Symposium

## Monte-Carlo simulation of ionization in Self-Induced Ion Plating (SIP)

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BRUXELLES

# Outline

- **Self-induced Ion Plating (SIP)**
- **SIP modelling**
  - **Magnetic model**
  - **Heat transfer model**
    - **Monte-Carlo method**
    - **Heat transfer model**
  - **Evaporation model**
- **Conclusions**

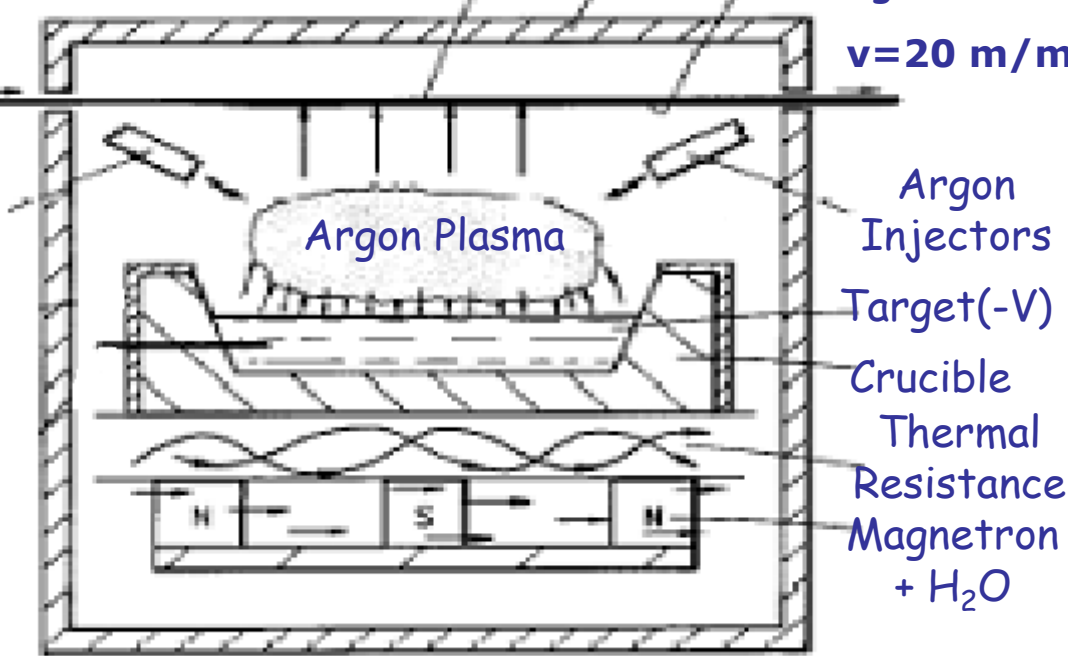
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# Self-induced ion plating

Substrate Enclosure Coating

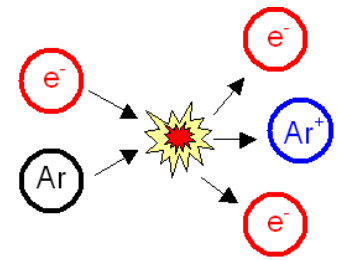
$v = 20 \text{ m/min}$



Tin target at low voltage : -900V



Electron displacement between the cathode and the anode



European Patent N°0780489

$p = 6,75 \cdot 10^{-4} \text{ Torr} = 0,09 \text{ Pa}$  ;

Target mean power density :  $21,5 \text{ W/cm}^2$

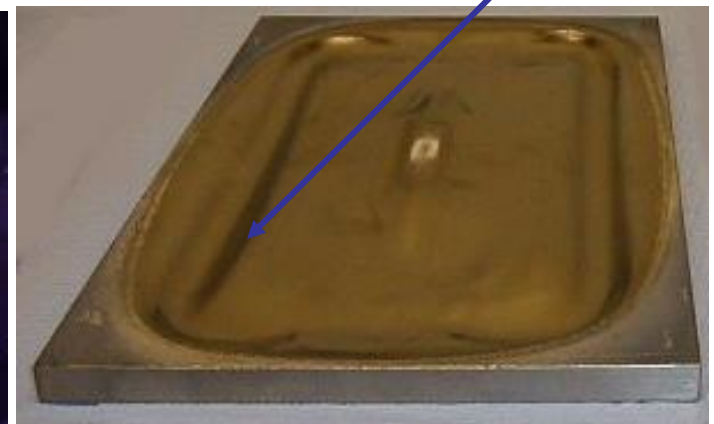
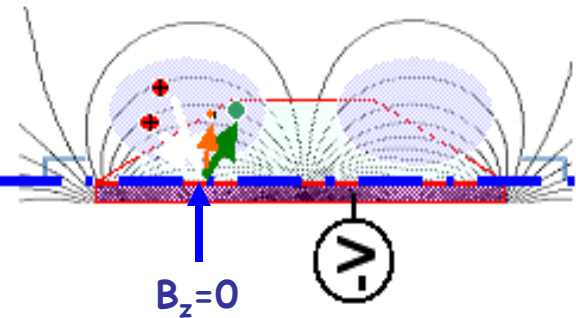
**Ar<sup>+</sup> and target at low voltage → Ion bombardment of the target**

**Thermal resistance under the crucible → Target evaporation**

⇒ Creation of Argon plasma

# Self-induced ion plating

- SIP = evaporation due to magnetron sputtering system
- Magnetron sputtering: Maximal **erosion** where the magnetic field is tangential to the target surface ( $B_z=0$ )



- SIP: Maximal **heat flux** where the magnetic field is tangential to the target surface ( $B_z=0$ )

→ Link between magnetic field and distribution of the ion bombardment heat flux : **Monte-Carlo method**

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# SIP modelling

## Magnetic Model



Magnetic field



Ion bombardment heat flux

## Heat transfer model



Temperature field at  
the target level

## Evaporation and coating model

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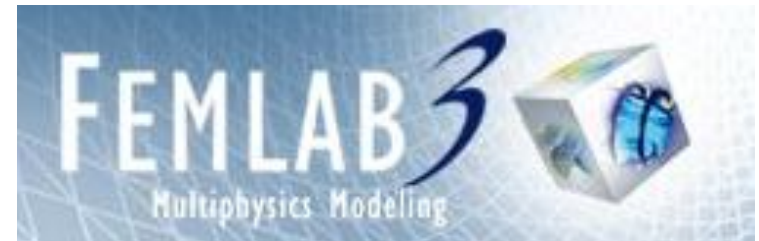


# Magnetic Model

**Magnetic field - FEMLAB® 3.0a**

**Magnetostatic laws**

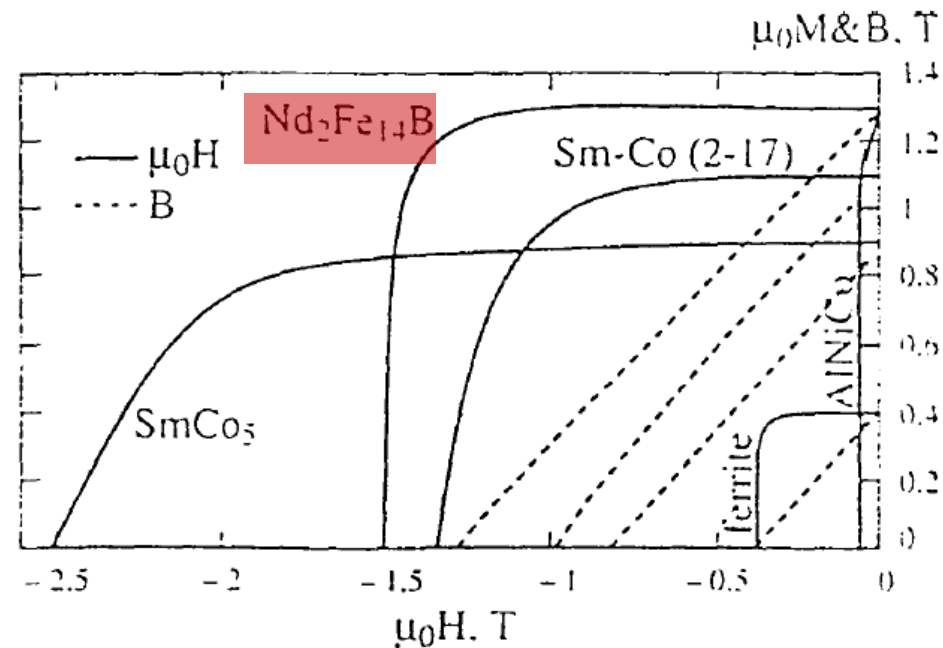
**Neodymium magnets (NdFeB)**



$$B = \mu_0(H + M); H = -\nabla V_m$$

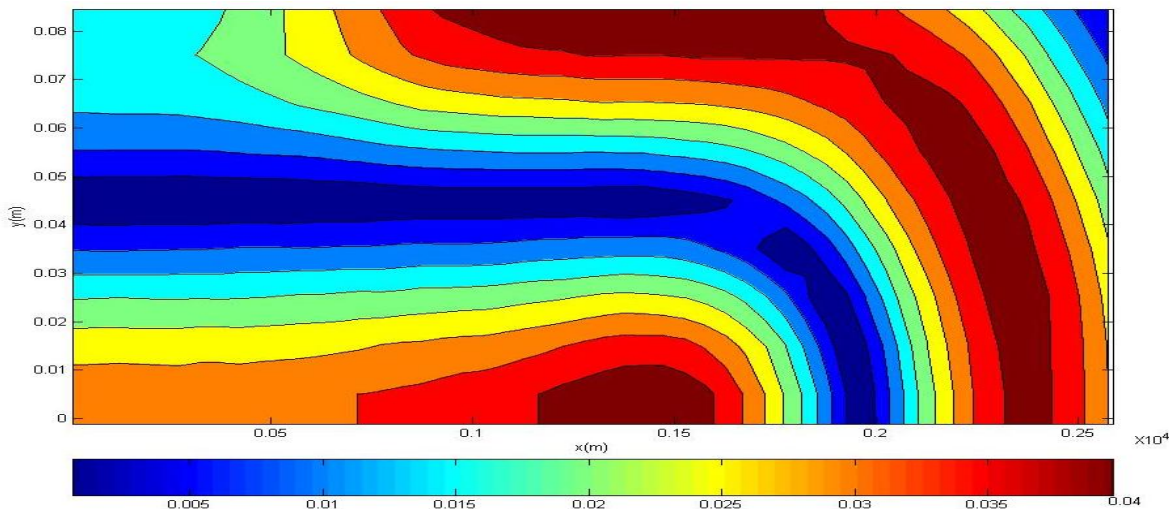
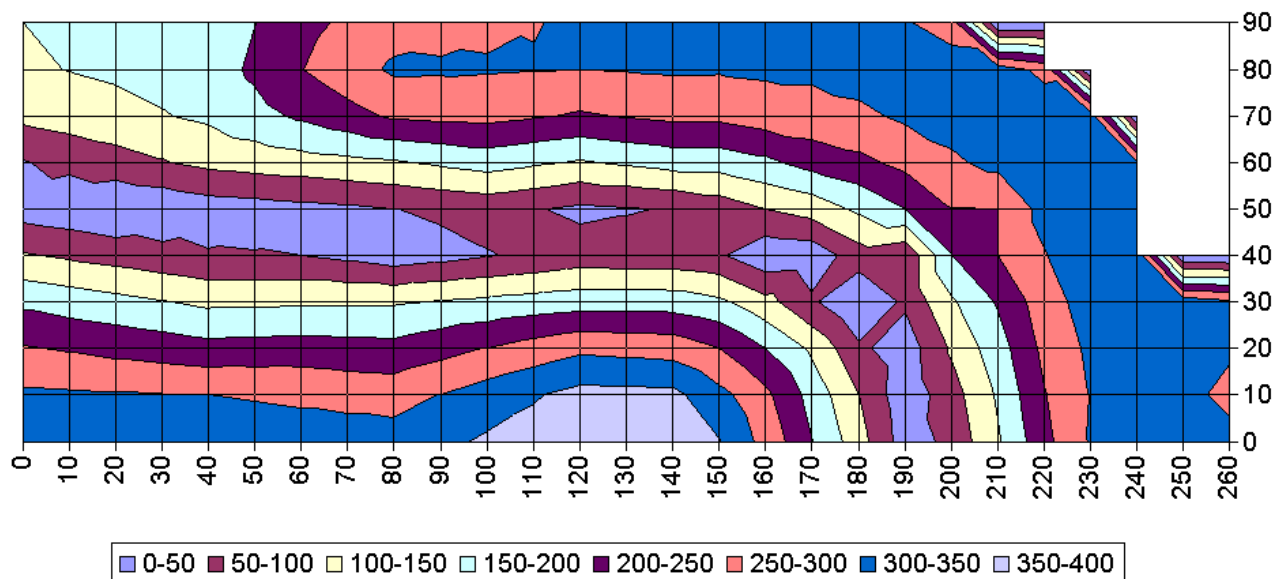
$$M = M_0 + (\mu_r - 1)H \Rightarrow B = \mu_0 \mu_r H + B_r$$

$$\nabla B = 0 \Rightarrow -\nabla(\mu_0 \mu_r \nabla V_m - B_r) = 0$$



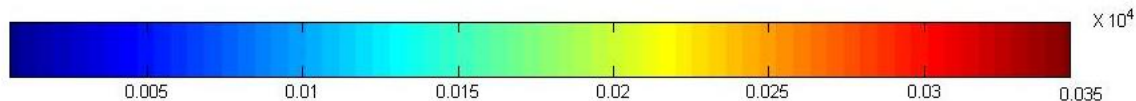
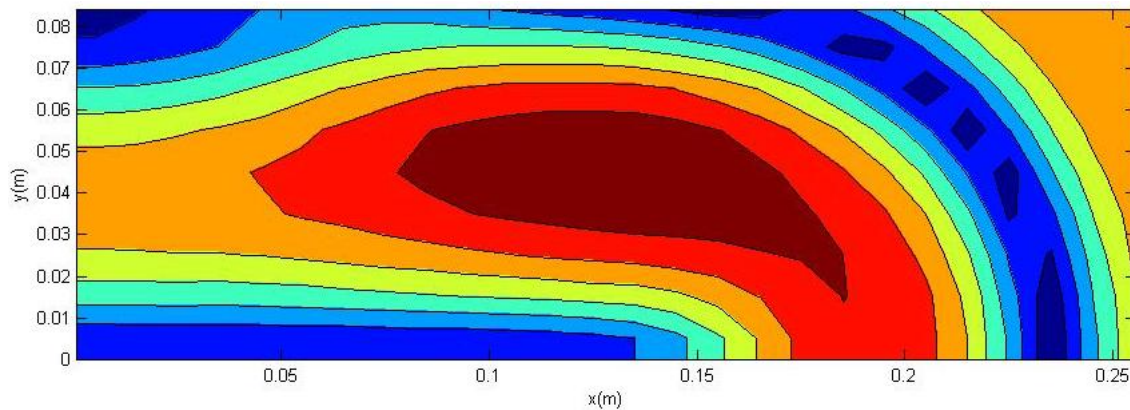
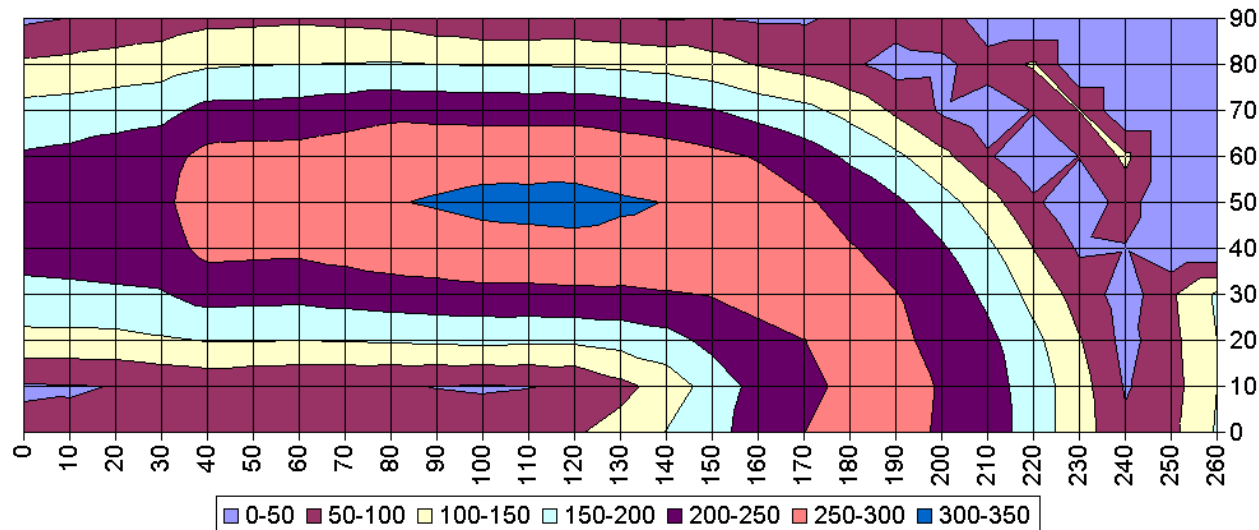
# Magnetic model

$|B_z|$  (Gauss)  
at the  
target level



# Magnetic model

$|B_{xy}|$  (Gauss)  
at the  
target level

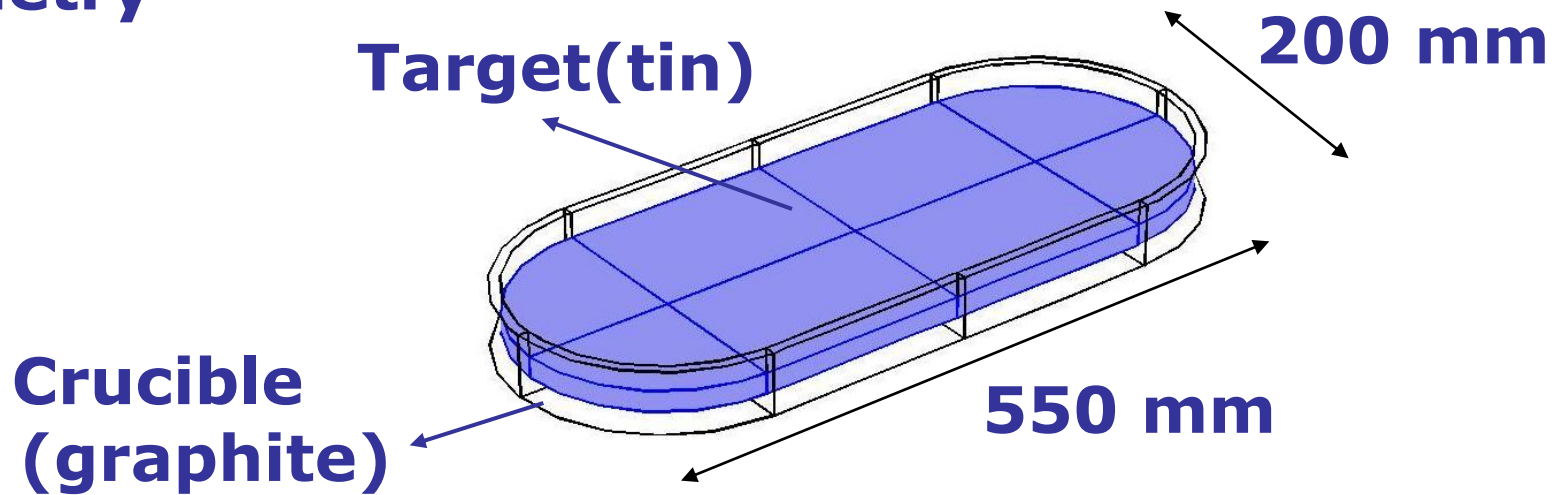


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# Heat transfer model

## Geometry



**Distribution of ion bombardment heat flux unknown**

→ **Monte-Carlo method**

→ **Distribution of bombardment heat flux =  
Distribution of ionisation points (ions not  
influenced by magnetic field)**

# Monte-Carlo method

## Step 1. Secondary electrons trajectories

$$\vec{F} = q(\vec{E} + \vec{v} \times \vec{B}) \Rightarrow \frac{\partial \vec{v}}{\partial t} = \frac{q}{m}(\vec{E} + \vec{v} \times \vec{B}) \quad (\text{Lorentz law})$$

## Numerical integration – Runge-Kutta (4th order)

## Magnetic field - Magnetostatic laws

## Electric field - Analytical model : Sheridan & Goree\*

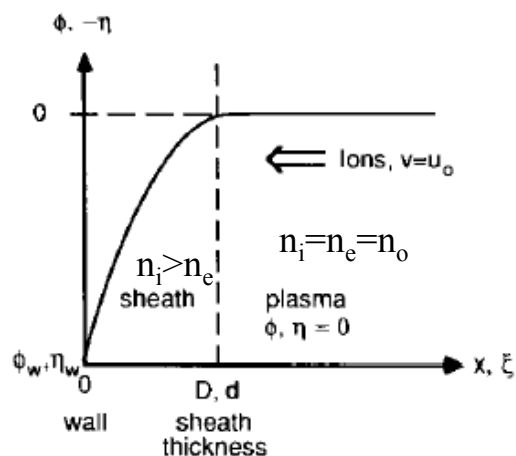
\*T. Sheridan, J. Goree, « Analytic expression for the electric potential in the plasma sheath », IEEE Transactions on plasma science 17(6) (1989) p. 884

### Sheath is source free → continuity

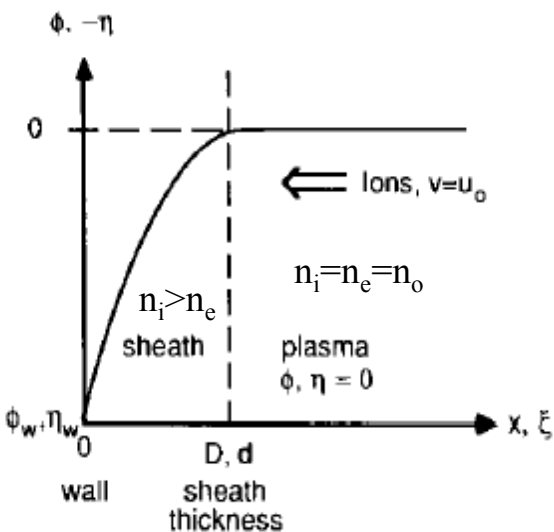
$$n_i v = n_0 v_0$$

### Sheath is collisionless → energy conservation

$$\frac{Mv^2}{2} = \frac{Mu_0^2}{2} + e\phi$$



# Monte-Carlo method



$$\rightarrow n_i = n_0 \left(1 - \frac{2e\phi}{Mu_0^2}\right)^{\frac{-1}{2}}$$

**Boltzmann relation**  $\rightarrow n_e = n_0 e^{\frac{e\phi}{kT_e}}$

**Poisson's relation**  $\rightarrow \frac{d^2\phi}{dx^2} = -\frac{e}{\epsilon_0} (n_i - n_e)$

$$\rightarrow \frac{d^2\phi}{dx^2} = -\frac{en_0}{\epsilon_0} \left[ \left(1 - \frac{2e\phi}{Mu_0^2}\right)^{\frac{-1}{2}} - e^{\frac{e\phi}{kT_e}} \right]$$

$$M_a = \frac{u_0}{\left(\frac{kT_e}{M}\right)^{\frac{1}{2}}}$$

**Dimensionless parameters :**  $\eta = \frac{e\phi}{kT_e}$ ;  $\xi = \frac{x}{\lambda_D}$ ;  $M_a = \frac{u_0}{\left(\frac{kT_e}{M}\right)^{\frac{1}{2}}}$

$\rightarrow$  if  $\eta_c > 10^4$  (Child's law)

$$\eta(\xi) \cong 1.36284(d - \xi)^{\frac{4}{3}}; d \cong 0.7928\eta_c^{\frac{3}{4}}$$

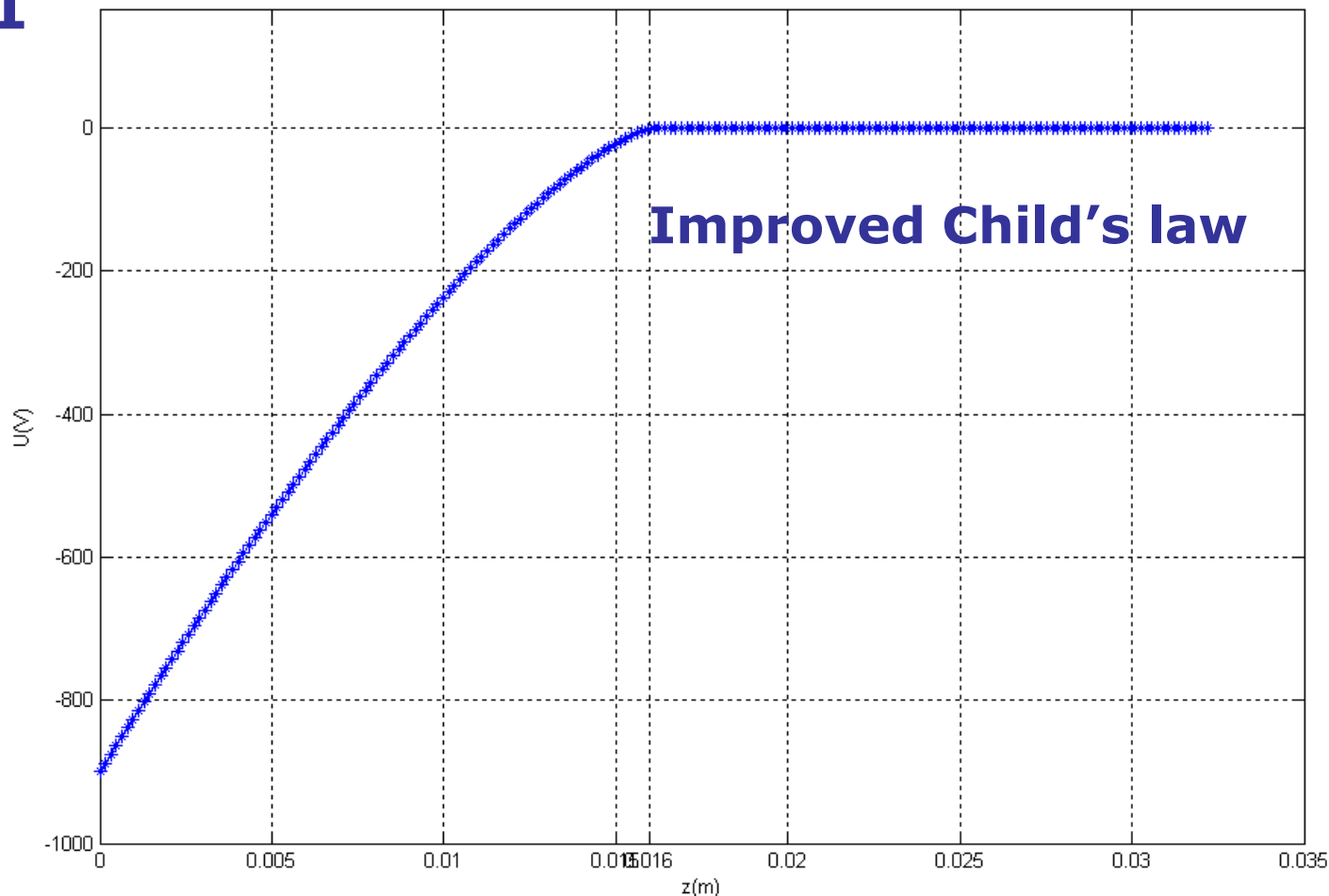
$\rightarrow$  if  $10^4 > \eta_c > 10$  (Improved Child's law)

$$\eta(\xi) = a(d - \xi)^b; a = \frac{\eta_c}{d^b}; d = \frac{\eta_c \eta_c}{\eta_c \eta_c - \eta_c^2}; b = \frac{1}{1 - \frac{\eta_c \eta_c}{\eta_c^2}}$$

# Monte-Carlo method

Improved Child's law for SIP :  $T_e = 4\text{eV}$ ;  $\Phi = -900\text{V}$ ;

$\eta_c = 261$

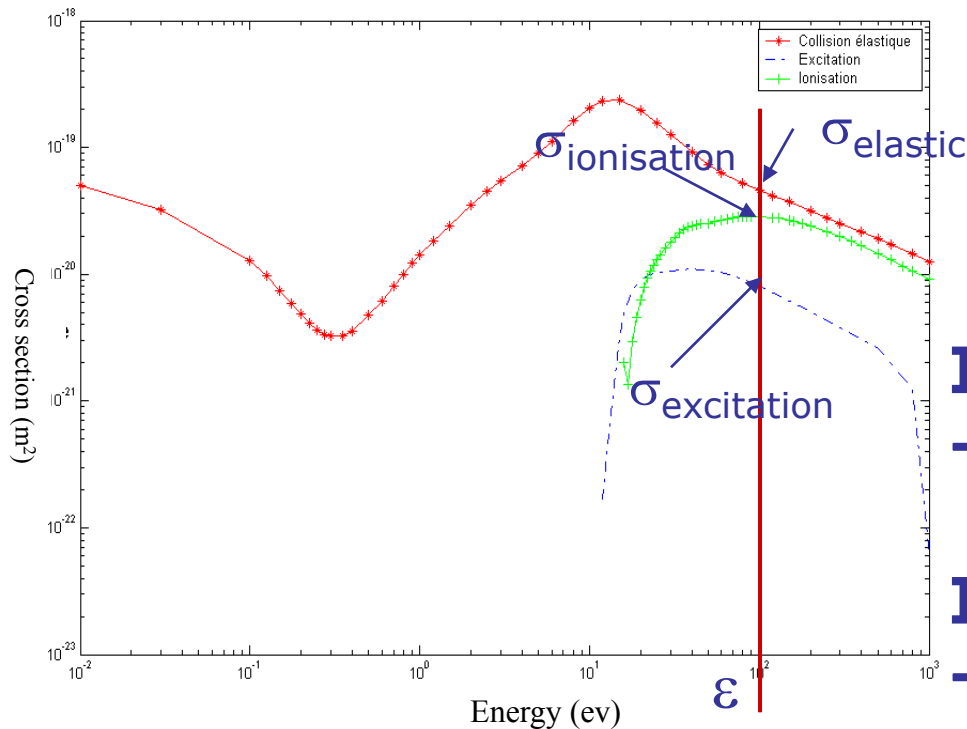


**Width of the sheath : 16 mm**



# Monte-Carlo method

## Step 2. At each time step : collision or not ?



### Probability of a collision

$$P = 1 - e^{-v\Delta t \sigma_{tot}(\epsilon)ng}$$

$$\sigma_{tot}(\epsilon) = \sigma_{elast}(\epsilon) + \sigma_{excit}(\epsilon) + \sigma_{ionis}(\epsilon)$$

If  $R > P \rightarrow$  No collision  
 $\rightarrow$  integration in  $t + \Delta t$

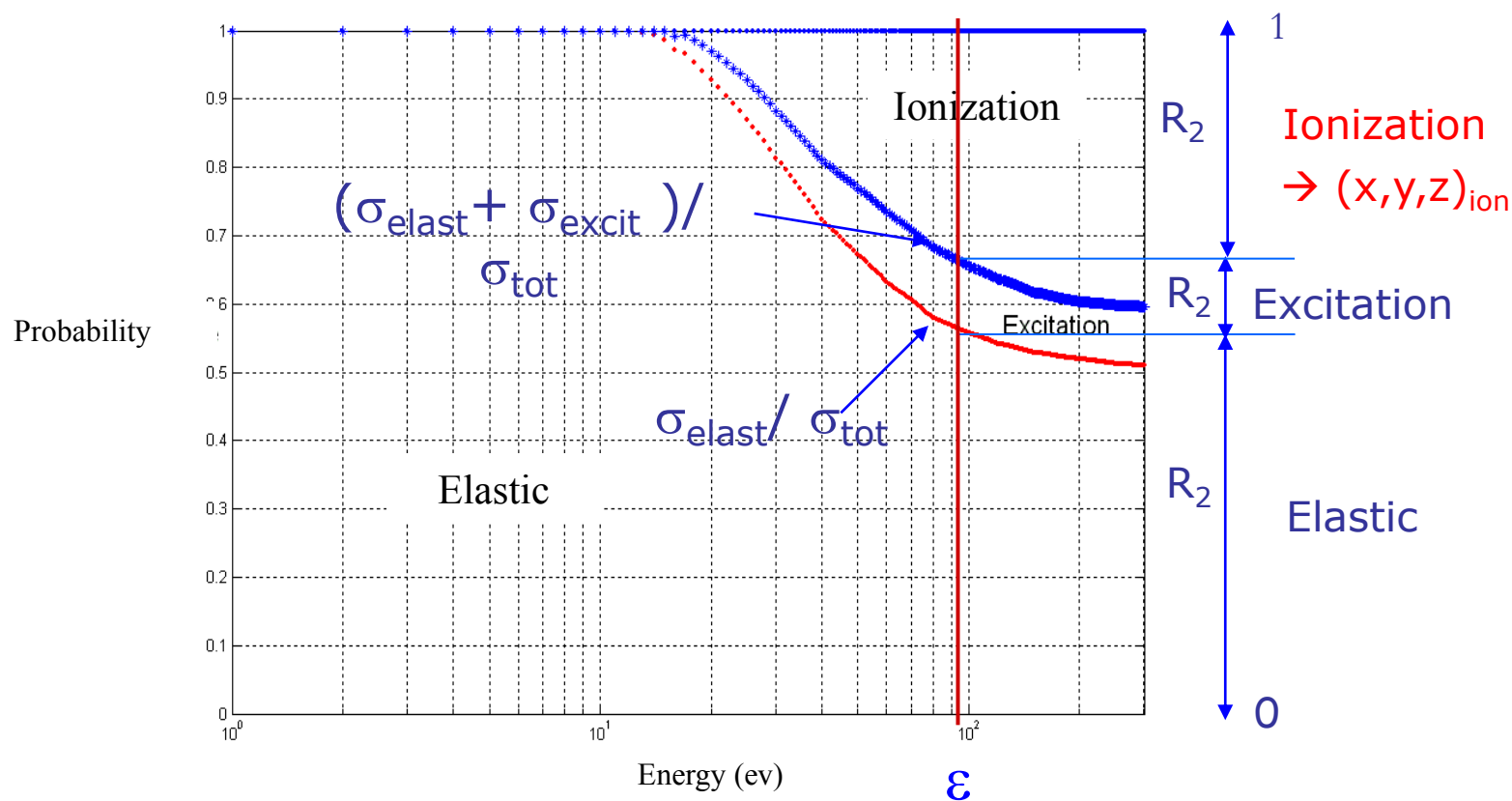
If  $R < P \rightarrow$  Collision  
 $\rightarrow$  Definition of the collision type

$R$ , random number uniformly distributed in the interval (0,1)

# Monte-Carlo method

## Step 3. Definition of the collision type and consequences on the incident electron orbit

$R_2$ , random number uniformly distributed in the interval (0,1)



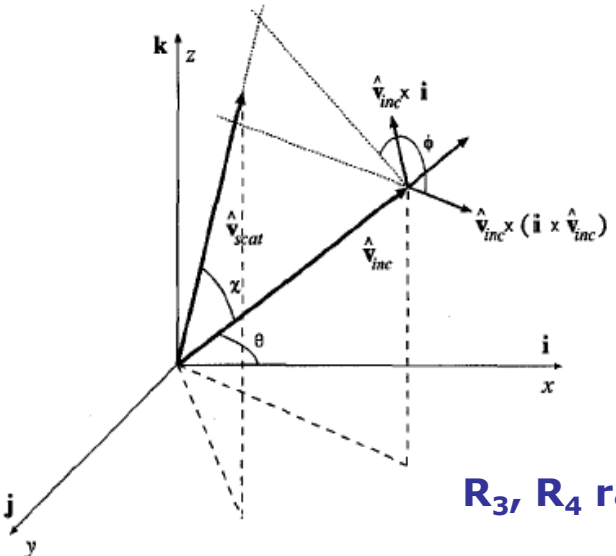
# Monte-Carlo method

## Consequences :

- Decreasing of the incident electron energy

Collision	Energy decreasing (eV)
Elastic	0
Excitation	11,6 (Excitation pot.)
Ionization	15,8 (Ionization pot. )

- Scattering of the incident electron velocity

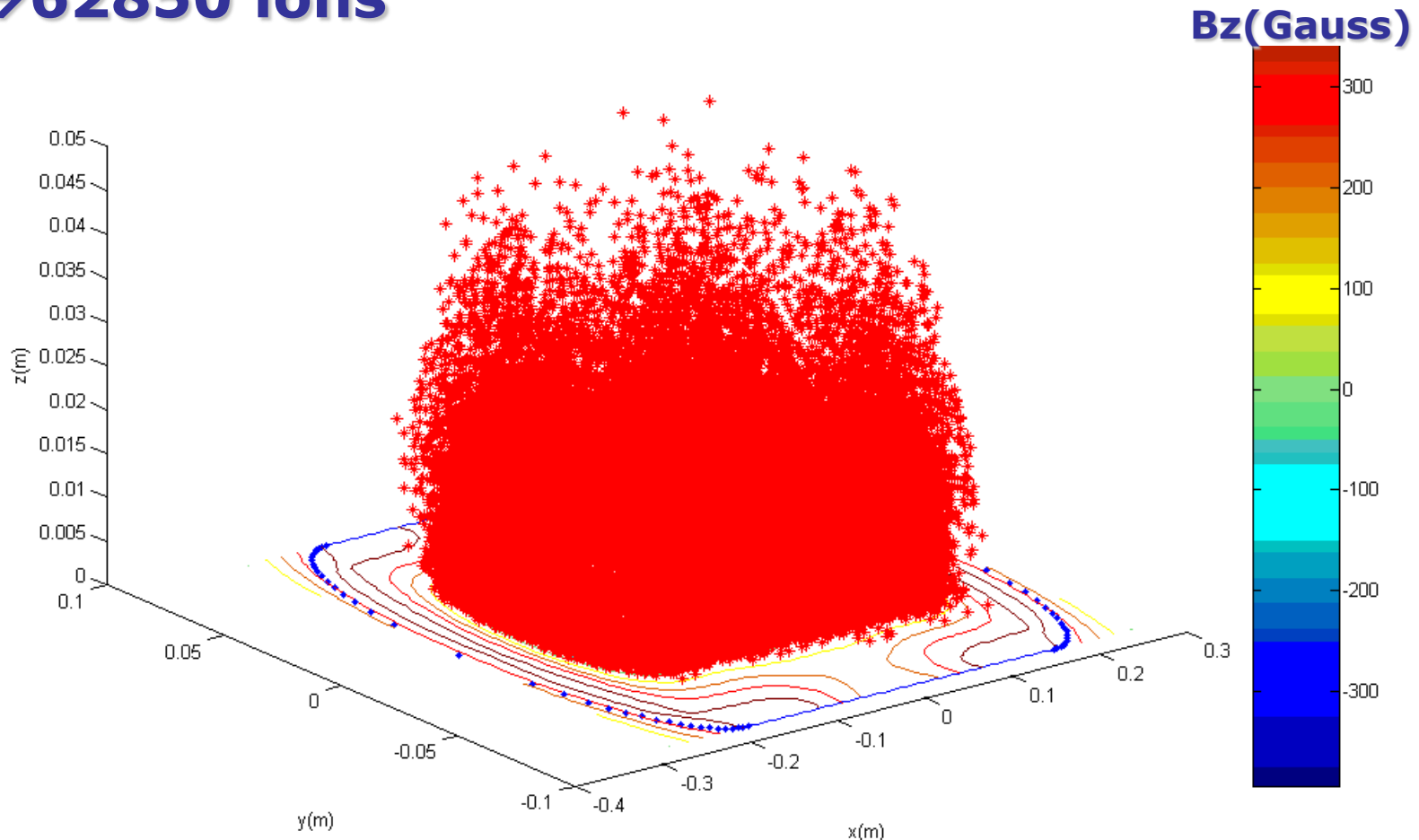


Scattering angles
$\chi = \frac{2 + \varepsilon - 2(1 + \varepsilon)^{R_3}}{\varepsilon}$ $\phi = 2\pi R_4$

$R_3, R_4$  random numbers uniformly distributed in the interval (0,1)

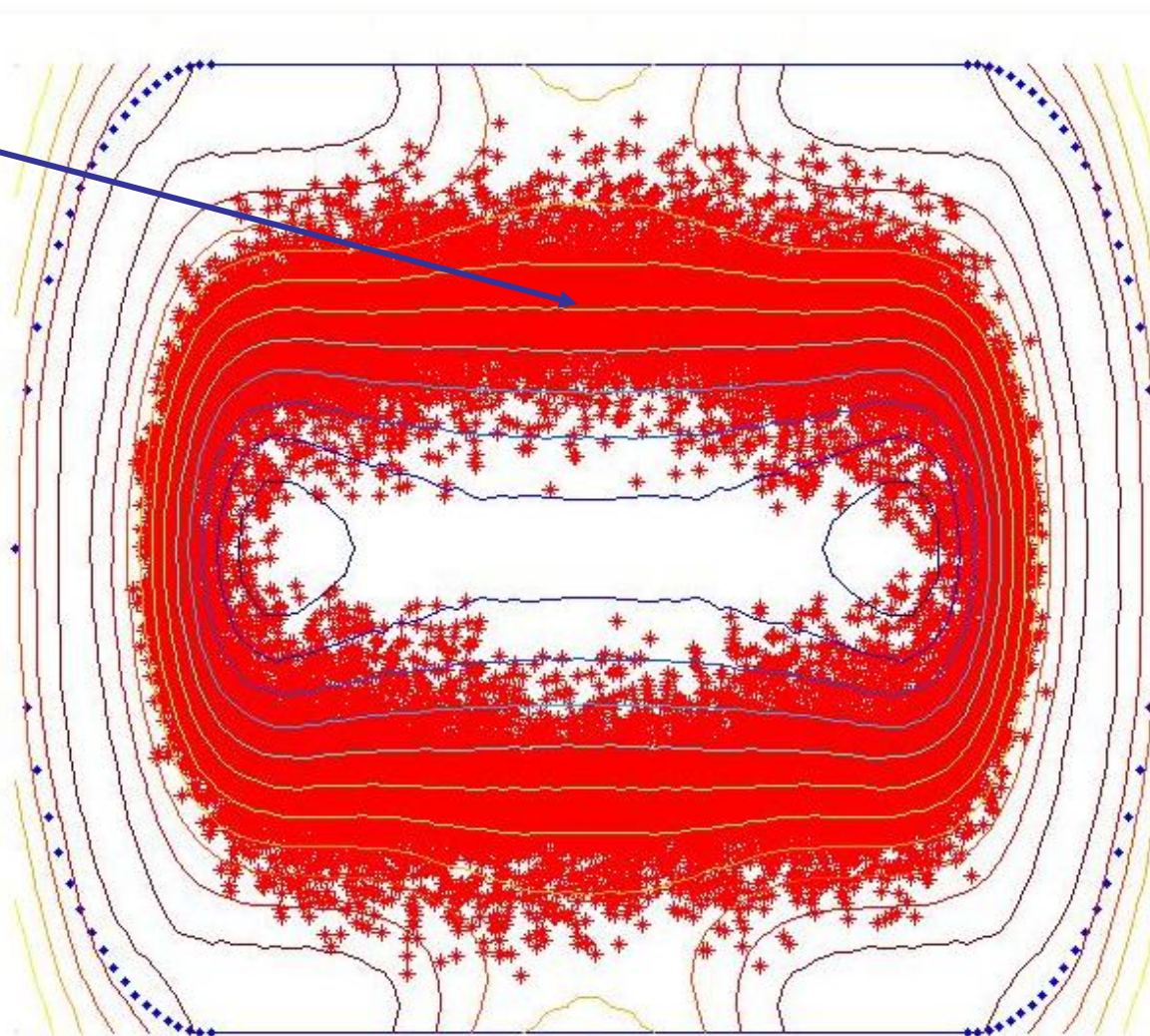
# Monte-Carlo - SIP

Distribution of ionization points : 34000 secondary electrons from the projections of ionization sites on target  $\rightarrow$  62850 ions

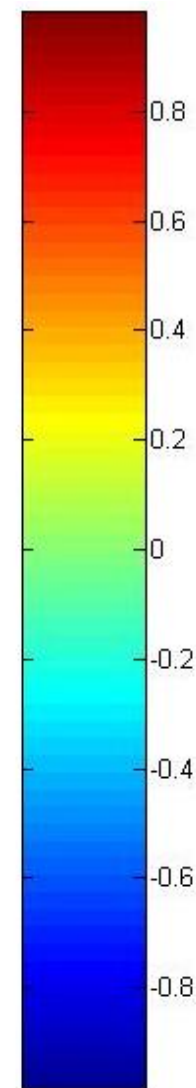


# Monte-Carlo - SIP

$B_z=0$



$B_z/B_{z_{max}}$

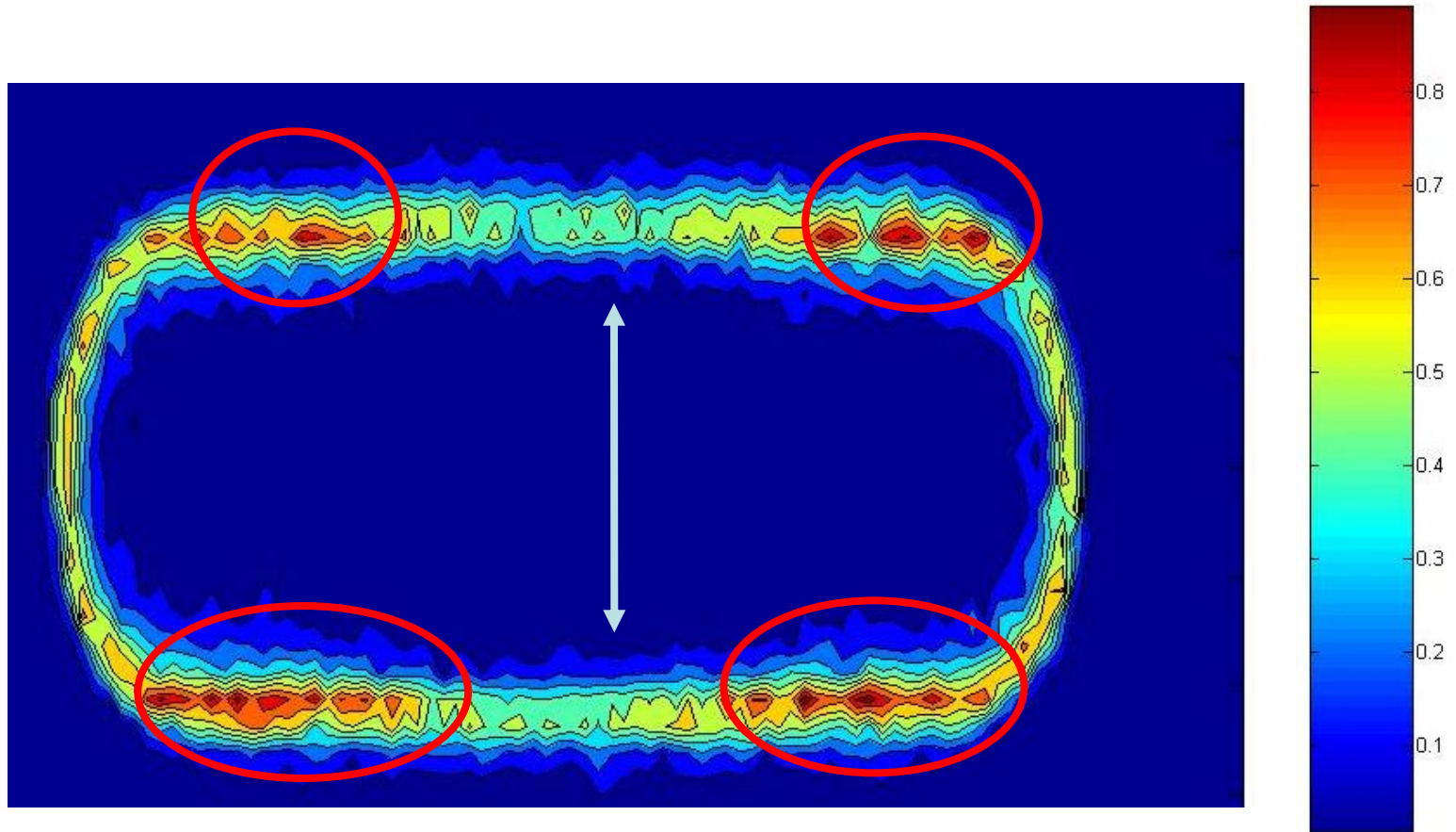


→ race-track around  $B_z=0$

# Monte-Carlo - SIP

## Normalised distribution of ionization points

→ Normalised distribution of ion bombardment heat flux



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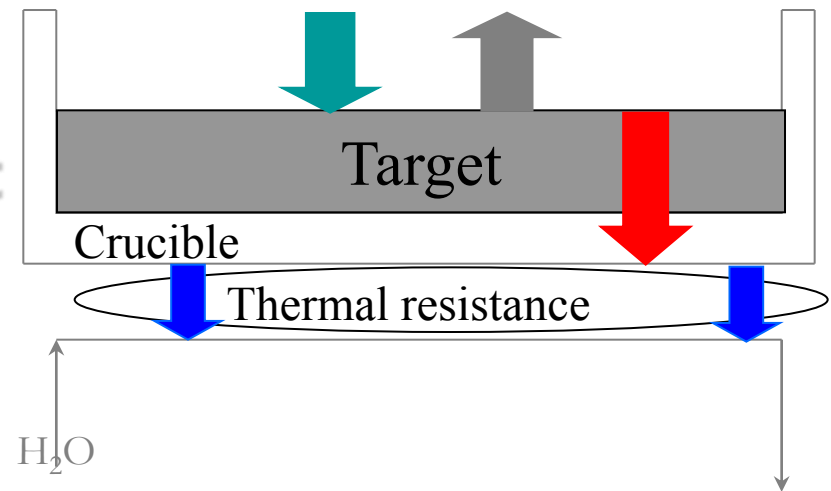
# Heat transfer model

FEMLAB® 3.0a

- **Conduction through target and crucible**

Lower crucible surface

- **Exchange between the lower crucible surface and the water flow through thermal resistance**



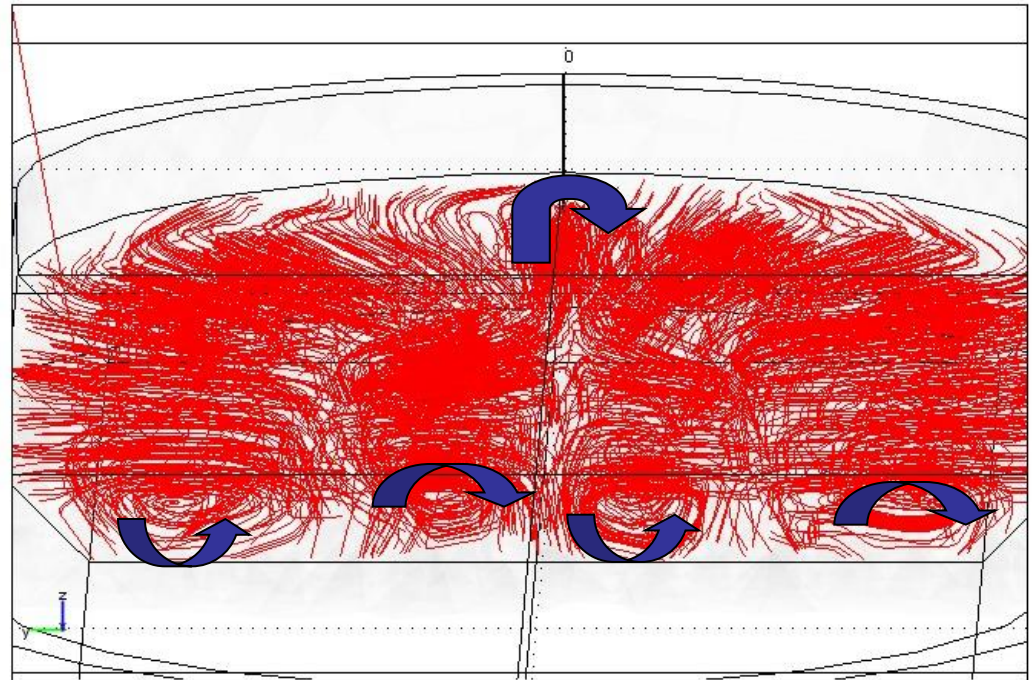
Target Level

- **Heat flux due to the ion bombardment (MC)**
- Heat flux lost because of target evaporation



# Heat transfer model

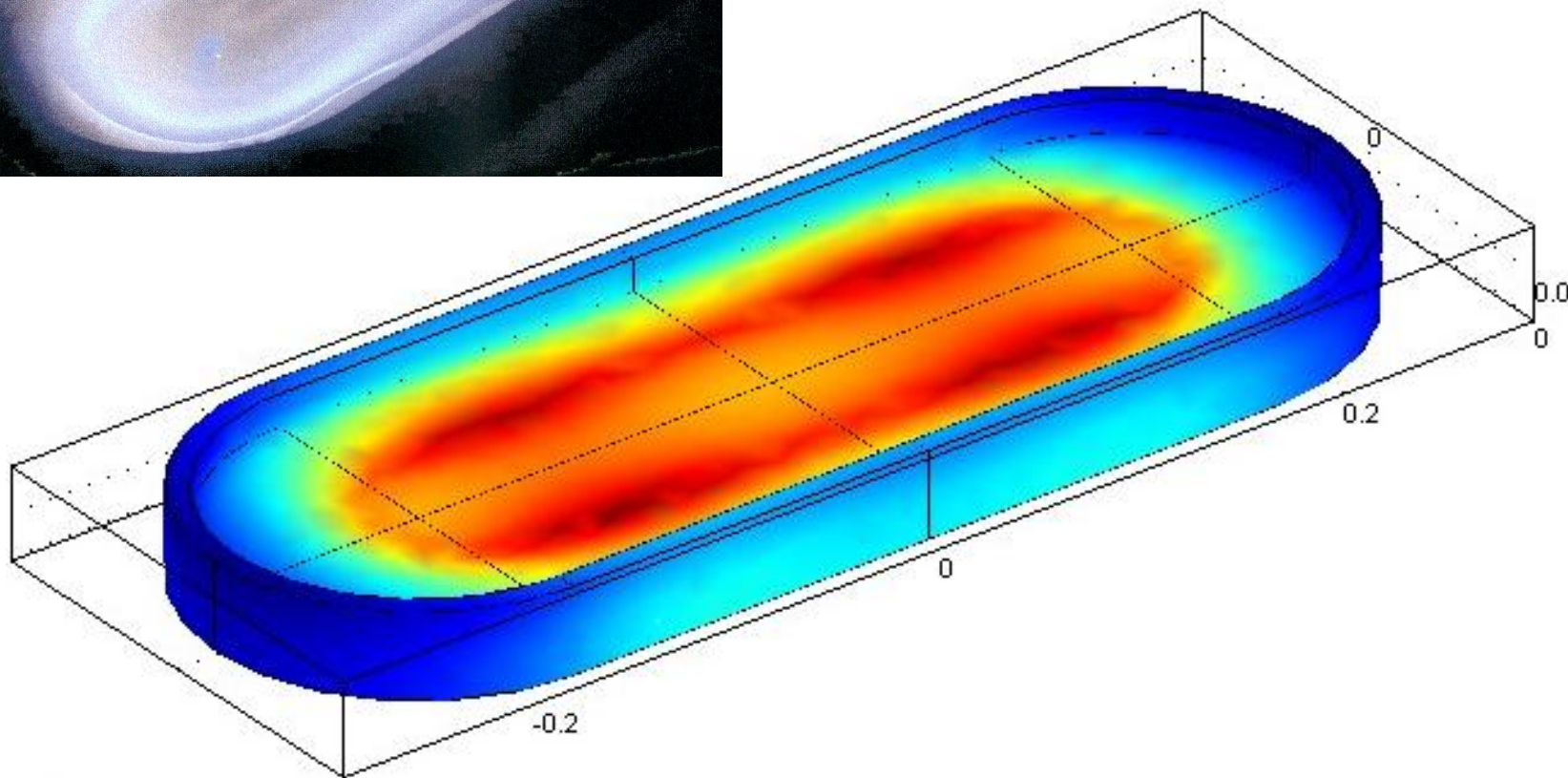
- Free convection in the liquid target due to variation of the density with temperature



- Radiative heat transfer between target, crucible, substrate and enclosure not taken into account accurately with Femlab

→ Need of improvement of the heat transfer model

# Heat transfer model



Max: 1817.025

1800

T(K)

1700

1600

1500

1400

1300

Min: 1227.953

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# Evaporation model

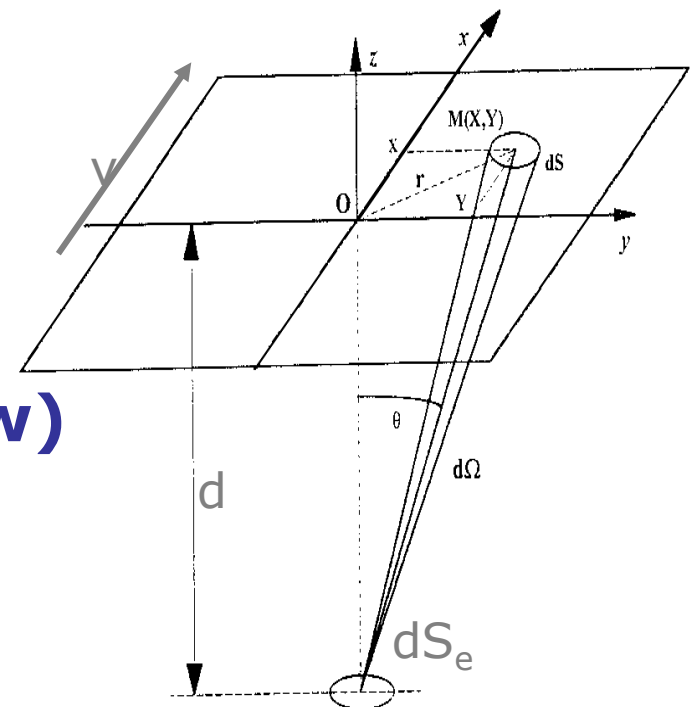
Temperature at the target level → thickness coating profile

Theory of evaporation from a point source

$$v_d = \sum_{i=1}^{i=nbr_{evap}} \frac{J_{ei} S_{ei} h^2}{\pi \rho [(x_d - x_{ei})^2 + (y_d - y_{ei})^2 + h^2]^2}$$

$$J_{ei} = 0,5834 \sqrt{\frac{M}{T}} p_v \quad \text{(Knudsen's law)}$$

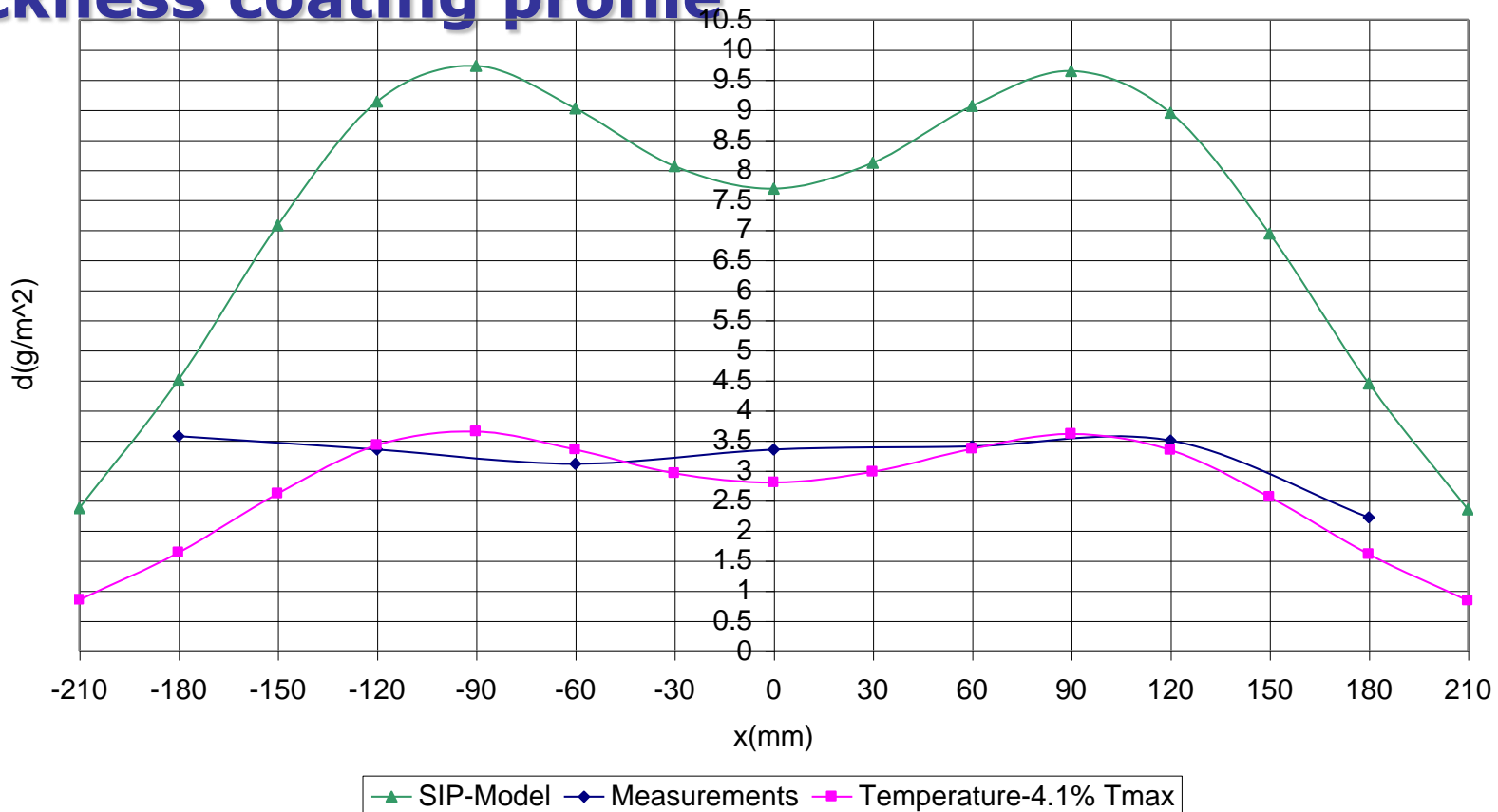
$$\text{Tin: } p_v = 0.0075 * 10^{10.268 - \frac{15332}{T}}$$



Displacement of the substrate taken into account

# Evaporation model

## Thickness coating profile



- Overestimation of the thickness measurements → radiation computation
- No symmetry of measurements → new tests planned
- System very sensitive to temperature variation

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# Conclusions

- **Modelling of the SIP**

**Magnetic model → Heat transfer model → Evaporation model**

- **Monte-Carlo method → ion bombardment heat flux distribution**

- **Improvements of the model :**

- **Computation of radiative heat transfer**
- **New tests are planned to validate model**