

Monte Carlo Simulations of the T_2 relaxation induced by Cubic Shaped Superparamagnetic Nanoparticles

Florent Fritsche¹, Gilles Rosolen², Alice De Corte², Bjorn Maes², Yves Gossuin¹, and Quoc Lam Vuong¹

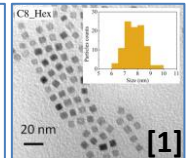
¹Biomedical Physics Unit, UMONS, 25 Avenue Maistriau, 7000 Mons, Belgium

²Micro-and Nanophotonic Materials Group, UMONS, 25 Avenue Maistriau, 7000 Mons, Belgium

Theoretical analysis of water protons transverse relaxation (T_2) induced by cubic-shaped superparamagnetic nanoparticles (Np) of magnetite, used as negative contrast agents in MRI, has been conducted with Monte Carlo simulations considering a high static magnetic field (B_0). The comparison between spherical and cubic-shaped nanoparticles, at equal volumes, revealed minor deviations in the transverse relaxation (T_2) within the Motional Average Regime [$d < 30\text{nm}$] whereas no deviation was observed for larger particles [$d > 30\text{nm}$].

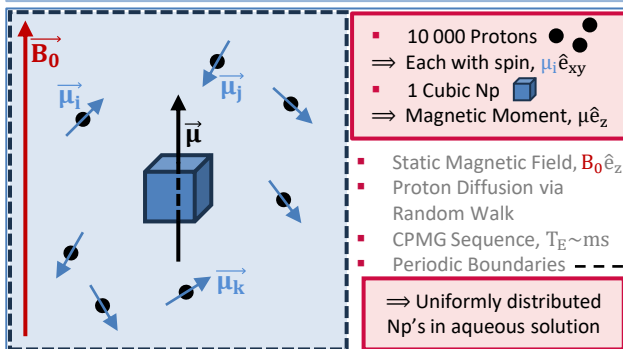
I. Introduction and research context

- For the last 20 years, there has been an ongoing interest in the synthesis and characterization of non-spherical nanoparticles. Several experimental studies, such as [1], using exotic particles reported an increase in the efficiency of shortening water relaxation times T_1 and T_2 .
- In contrast, very few studies have confronted experimental results to simulations or theory. Our goal is to study the influence of non-spherical shape induced relaxation by magnetite nanoparticles using Monte Carlo simulations.



II. Methodology

II. a. Simulation Setup [2]

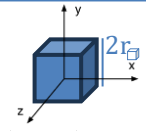


II. b. Simulations

Analytical Magnetic Field for a cubic particle [3]

$$\begin{cases} x_k = x + (-1)^k r_{\square} \\ y_l = y + (-1)^l r_{\square} \\ z_m = z + (-1)^m r_{\square} \end{cases}, \quad r_{klm} = \sqrt{x_k^2 + y_l^2 + z_m^2}$$

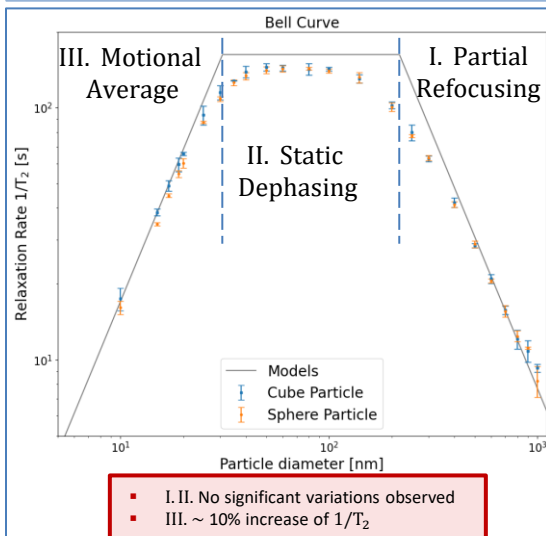
$$\Rightarrow B_z = \left(\frac{\mu_0 M_0}{4\pi} \right) \sum_{k,l,m=0}^1 (-1)^{k+l+m} \arctg \left(\frac{x_k y_l}{z_m r_{klm}} \right)$$



Temporal evolution of proton spin $\phi(\vec{r}(t))$:

- $\vec{\mu}_{i\perp}(t) = \mu_{i\perp}^0 (\cos(\phi), \sin(\phi))$ $\vec{\mu}_{i\parallel}(t) = \mu_{i\parallel}^0$
- $\Delta\phi(\vec{r}(t)) = \gamma B_z(\vec{r}) \Delta t \Rightarrow$ Larmor precession $\omega_0(t) = -\gamma B_z(\vec{r})$
- $|\vec{\mu}_{i\perp}(t)| \propto \exp(-t/T_2)$

III. T_2 - Comparizon between spherical and cubic contrast agents



I. Partial Refocusing (PR) Regime, [$d > 200\text{nm}$]

- Protons diffusing close to the Np are almost immediately lost from the signal due to high magnetic field gradients.
- \Rightarrow Relaxation arises mainly from proton diffusion far from the Np.
- As the distance grows, the magnetic field of a cubic shaped Np converges to the magnetic field of a spherical Np.
- \Rightarrow Proton spin dephasing in the PR is the **same for any shape**.

- It has been shown in [4] that $[1/T_2 \propto \sigma^2]$, where $[\sigma^2 = \frac{1}{V} \int |\nabla B_z|^2 dV]$. Study of cubic shape Np σ_{\square}^2 showed a steep convergence as the distance from its center grows to the spherical Np σ_{\circ}^2 ; confirming our interpretation.

II. Static Dephasing Regime, [$30\text{nm} < d < 200\text{nm}$]

- The statistical distribution of the magnetic field $[P(B_z)]$ converges rapidly as the distance grows to the dipole model. The lack of variation in relaxation can then be explained by the relation $[1/T_2 \propto p(B_z)]$ [5].

III. Motional Average Regime, [$d < 30\text{nm}$]

- Protons diffusing near the Np can perceive the contrast agent shape due to the small magnetic field intensity of the Np.
- Study of $[\langle B_z^2 \rangle]$ indicated a $\sim 20\%$ variation with the dipolar model which can be linked to our results via $[\langle B_z^2 \rangle \propto 1/T_2]$ by the Redfield theory.

IV. Summary and Future Directions

- Monte Carlo simulations demonstrate that the contrast agent shape has no impact on T_2 within the static diffusing and partial refocusing regimes. However, a 10% increase is observed in the Motional Average regime [$d < 30\text{nm}$].
- Future studies will focus on the agglomeration of exotic-shaped Np (stars, cylinders, ...) and their impact on T_2 in the MAR.

[1] Basini, M., Masciuffari, A., et al., "Local spin dynamics of IO magnetic Np dispersed in different solvents with variable size and shape: A 1H NMR study". *Chemical Physics*, 2017, Vol. 146, Issue 3. <https://doi.org/10.1063/1.4973979>

[2] Vuong, Q. L., Gillis, P., Roch, A., & Gossuin, Y. "Magnetic resonance relaxation induced by superparamagnetic particles used as contrast agents in magnetic resonance imaging: a theoretical review". *WIREs Nanomedicine and Nanobiotechnology*, 2017, Vol. 9, Issue 6. doi.org/10.1002/wnan.1468

[3] Engel-Herbert, R., & Hesjedal, T. "Calculation of the Magnetic Stray Field of a Uniaxial Magnetic Domain". *Journal of Applied Physics*, 2005, Vol. 97, Issue 7. doi.org/10.1063/1.1883308

[4] Majumdar, S., & Gore, J. C. "Studies of diffusion in random fields produced by variations in susceptibility". *Journal of Magnetic Resonance*, 1987, Vol. 78, Issue 1, pp. 41–55. [doi.org/10.1016/0022-2364\(88\)90155-2](https://doi.org/10.1016/0022-2364(88)90155-2)

[5] Brown, R. J. S. "Distribution of Fields from Randomly Placed Dipoles: Free-Precession Signal Decay as Result of Magnetic Grains". *Physical Review*, 1961, Vol. 121, Issue 5, pp. 1379–1382. doi.org/10.1103/physrev.121.1379