

Exciting Fano resonances in structured hyperbolic metamaterials

Fabio Vaianella^{*}, Bjorn Maes

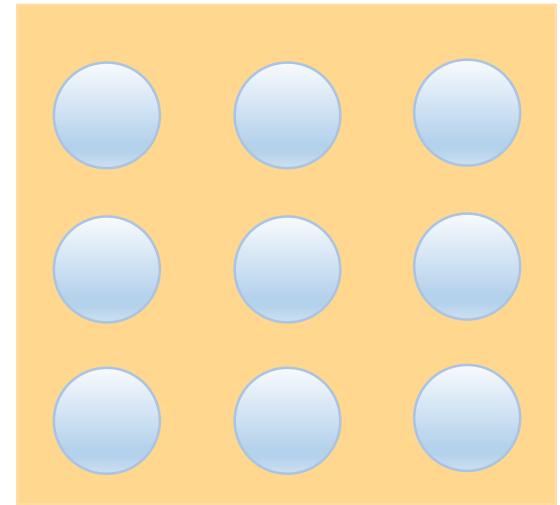
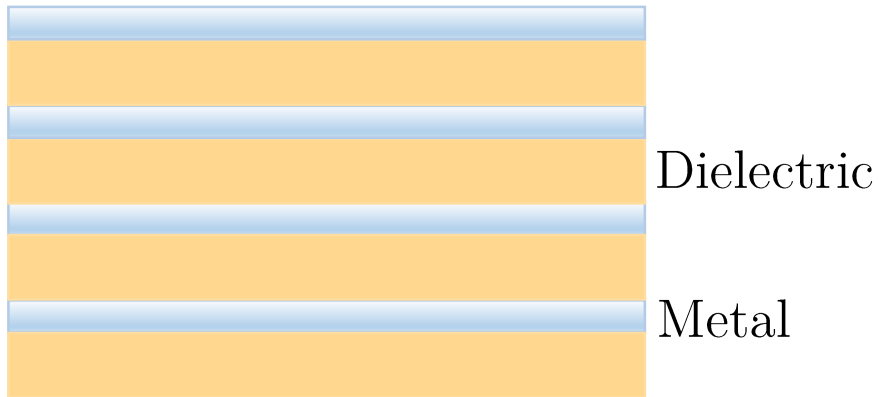
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- Introduction to hyperbolic metamaterials (HMMs)
- Some properties
- Hyperbolic cavities with Fano resonances

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Anisotropic media



Standard effective medium theory (Bruggeman):

$$\epsilon_{\parallel} = \begin{bmatrix} \epsilon_{\parallel} & 0 & 0 \\ 0 & \epsilon_{\parallel} & 0 \\ 0 & 0 & \epsilon_{\perp} \end{bmatrix}$$

$$\epsilon_{\parallel} = f \epsilon_m + (1 - f) \epsilon_d$$

$$\epsilon_{\perp} = \frac{\epsilon_m \epsilon_d}{\epsilon_m (1 - f) + \epsilon_d f}$$

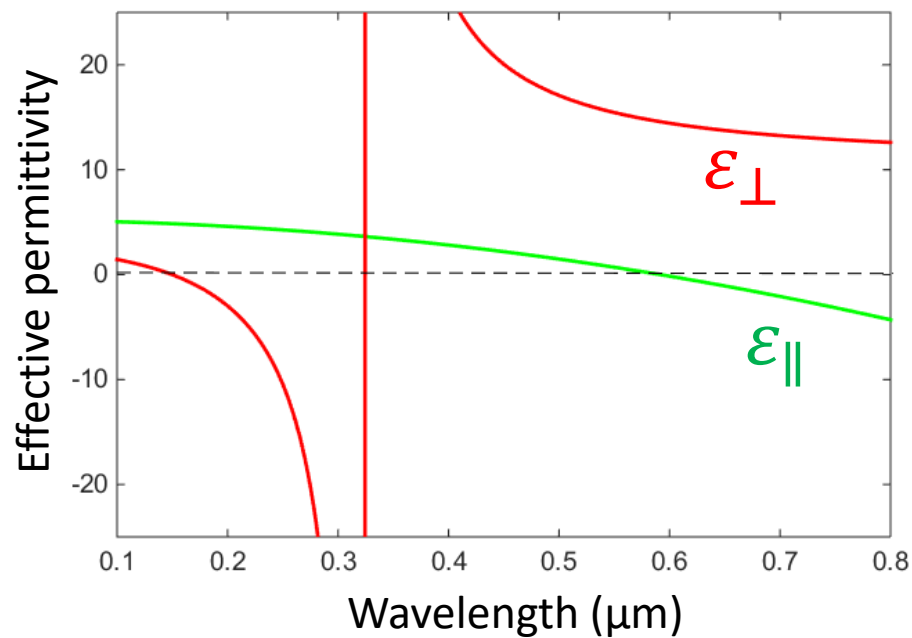
Metal fill factor ←

$$\frac{k_{\parallel}^2}{\epsilon_{\perp}} + \frac{k_{\perp}^2}{\epsilon_{\parallel}} = k_0^2$$

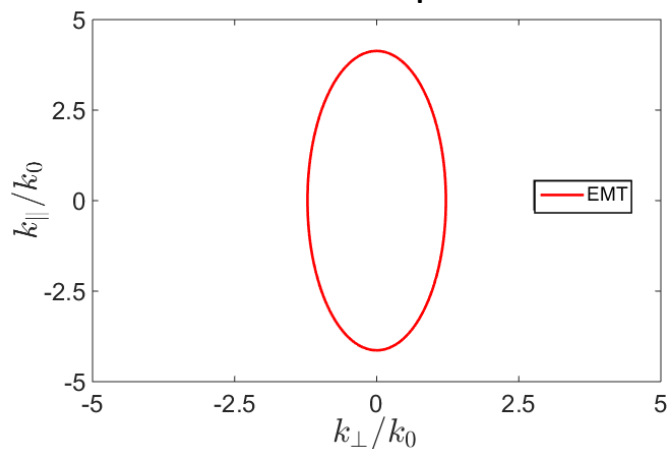
TM or p-polarization

Example with Ag and TiO₂

$f = 1/3$
 $d_{\text{Ag}} = 10 \text{ nm}$
 $d_{\text{TiO}_2} = 20 \text{ nm}$

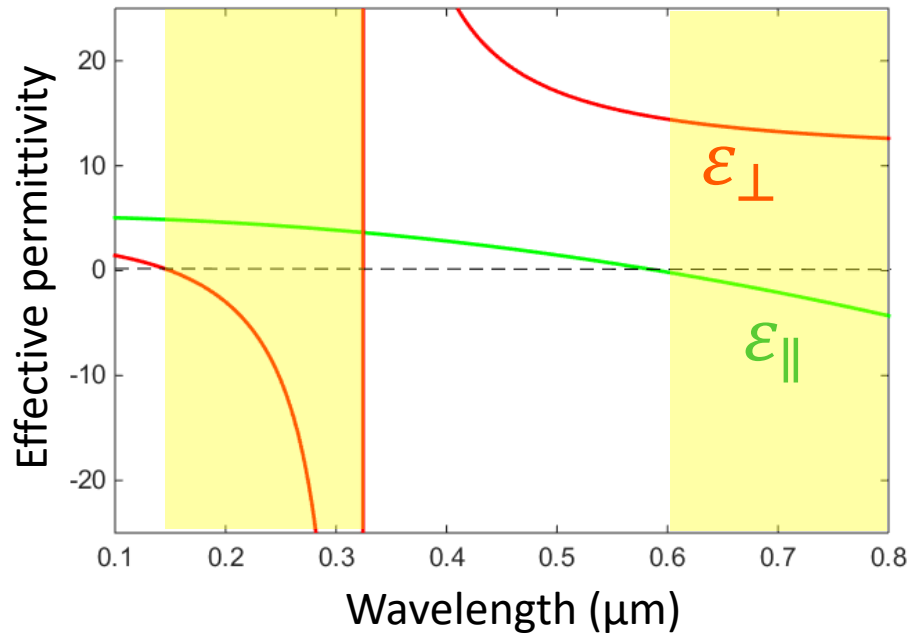


$\lambda = 500 \text{ nm}$ - elliptic



Example with Ag and TiO₂

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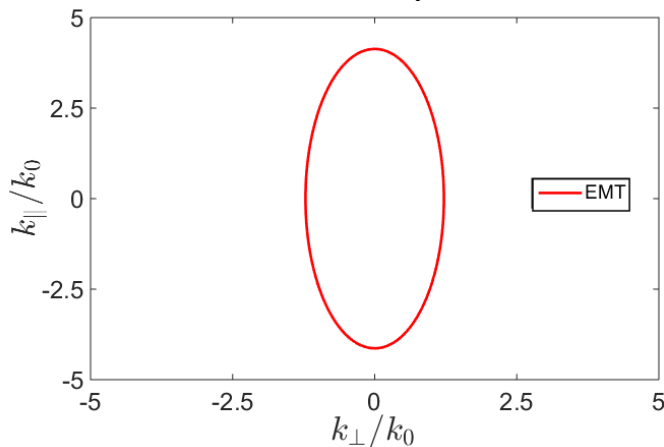


$\epsilon_{\parallel} \cdot \epsilon_{\perp} < 0$ possible

$$\frac{k_{\parallel}^2}{\epsilon_{\perp}} + \frac{k_{\perp}^2}{\epsilon_{\parallel}} = \frac{\omega^2}{c^2}$$

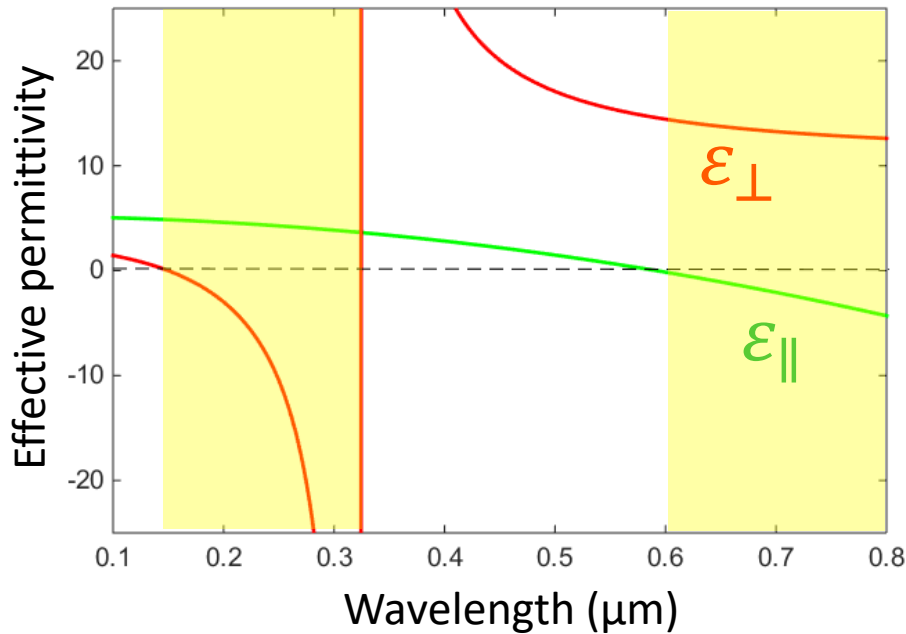
Hyperbolic isofrequency curve!

$\lambda = 500 \text{ nm}$ - elliptic



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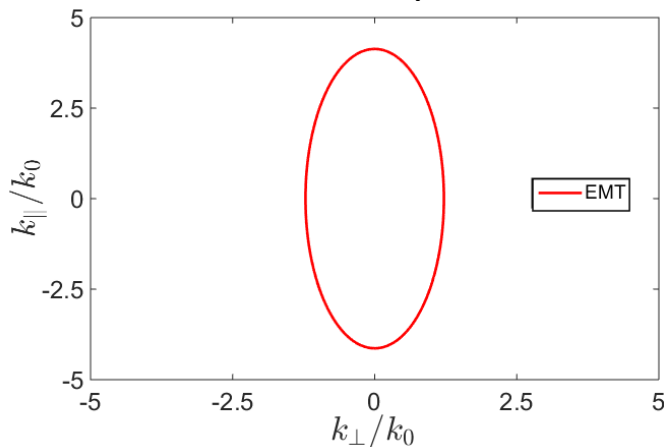


$\epsilon_{\parallel} \cdot \epsilon_{\perp} < 0$ possible

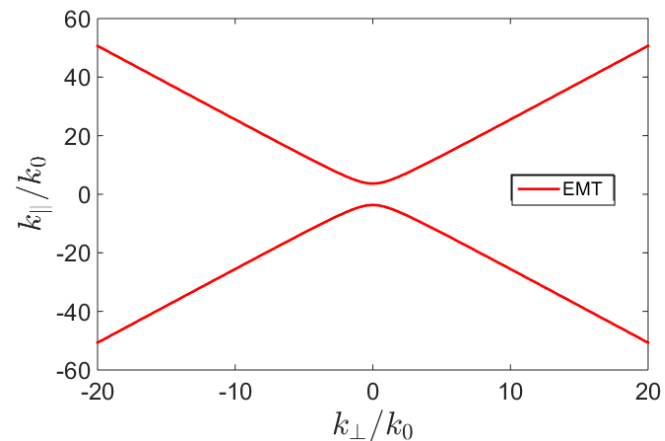
$$\frac{k_{\parallel}^2}{\epsilon_{\perp}} + \frac{k_{\perp}^2}{\epsilon_{\parallel}} = \frac{\omega^2}{c^2}$$

Hyperbolic isofrequency curve!

$\lambda = 500 \text{ nm}$ - elliptic



$\lambda = 700 \text{ nm}$ - hyperbolic



Limits of EMT

$$\frac{k_{\parallel}^2}{\varepsilon_{\perp}} + \frac{k_{\perp}^2}{\varepsilon_{\parallel}} = \frac{\omega^2}{c^2}$$

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Origin of hyperbolic properties: plasmonic
→ Nonlocality

Limits of effective medium theory

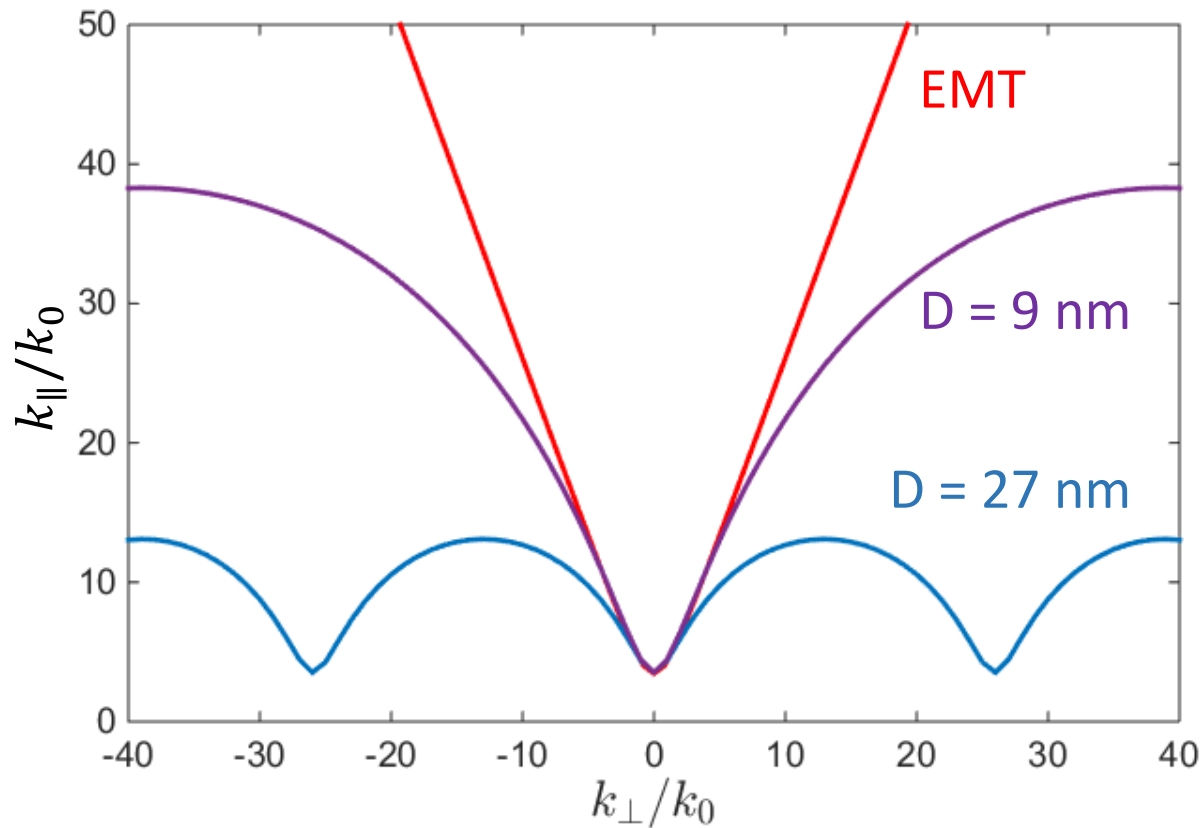
$$\cos(k_y D) = \frac{(\kappa_d \varepsilon_m + \kappa_m \varepsilon_d)^2}{4\kappa_d \kappa_m \varepsilon_d \varepsilon_m} \cosh(\kappa_d d_d + \kappa_m d_m) - \frac{(\kappa_d \varepsilon_m - \kappa_m \varepsilon_d)^2}{4\kappa_d \kappa_m \varepsilon_d \varepsilon_m} \cosh(\kappa_d d_d - \kappa_m d_m)$$

$$\kappa_{m,d} = \sqrt{k_x^2 - \varepsilon_{m,d} k_0^2}$$

Limits of effective medium theory

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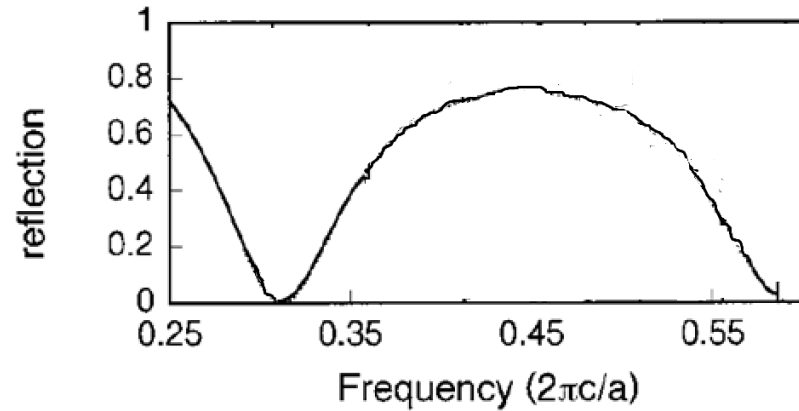
Limited inside
Brillouin zone:

$$\frac{\pi}{D}$$

Standard effective medium approach (EMT) not valid in many case

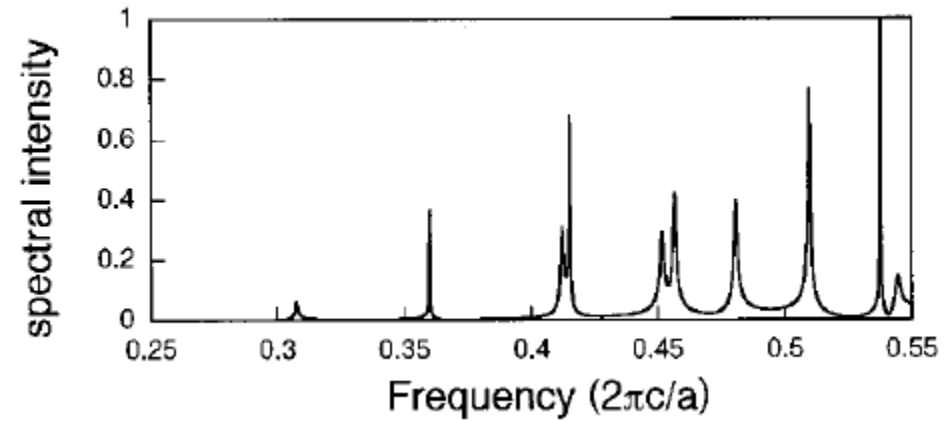
Fano resonances

Slowly varying background



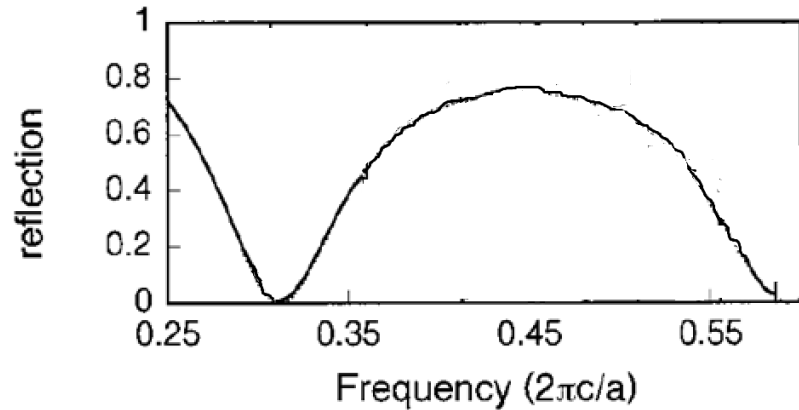
+

Narrow resonances



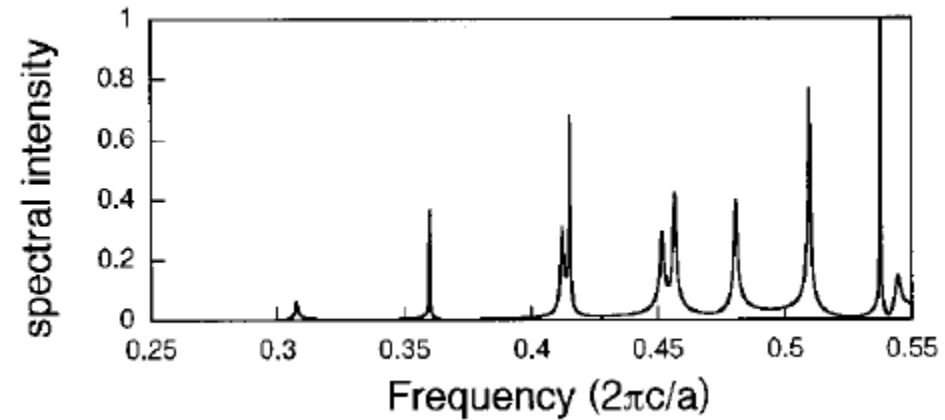
Fano resonances

Slowly varying background



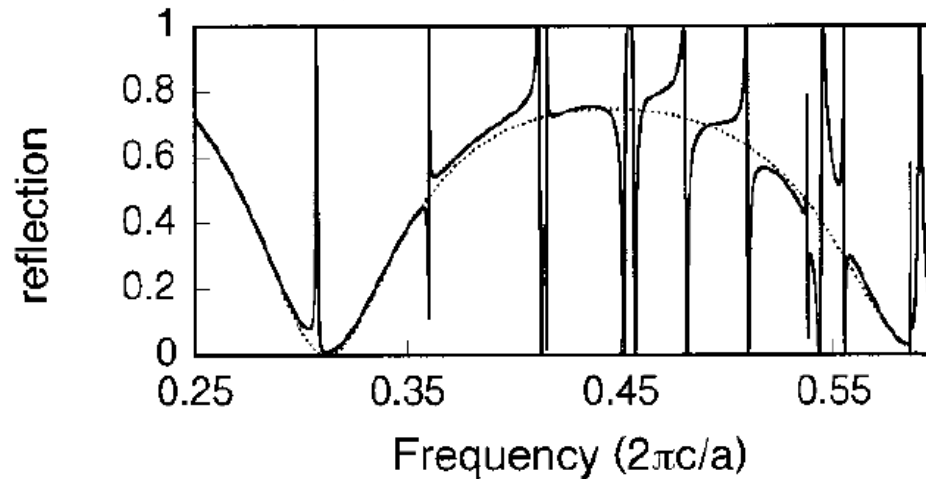
+

Narrow resonances



=

Asymmetric Fano resonances



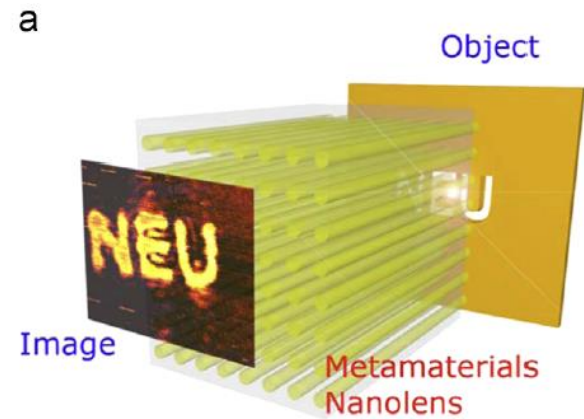
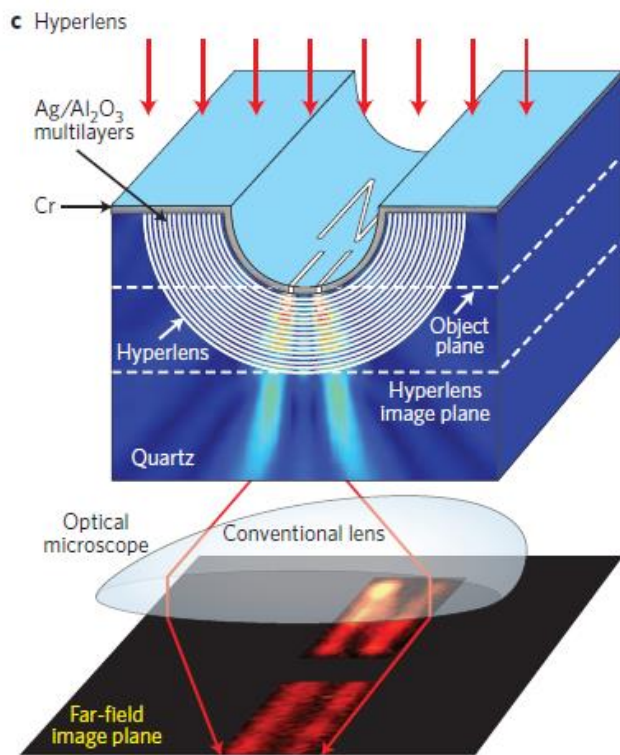
S. Fan and J.D.
Joannopoulos, Phys. Rev.
B, vol. 65, 235112. (2002)

- Introduction to hyperbolic metamaterials (HMMs)
- **Some properties**
- Hyperbolic cavities with Fano resonances

High-k propagating waves

High-k waves can propagate inside HMM \rightarrow Possibility to overcome diffraction limit

Application: hyperlens

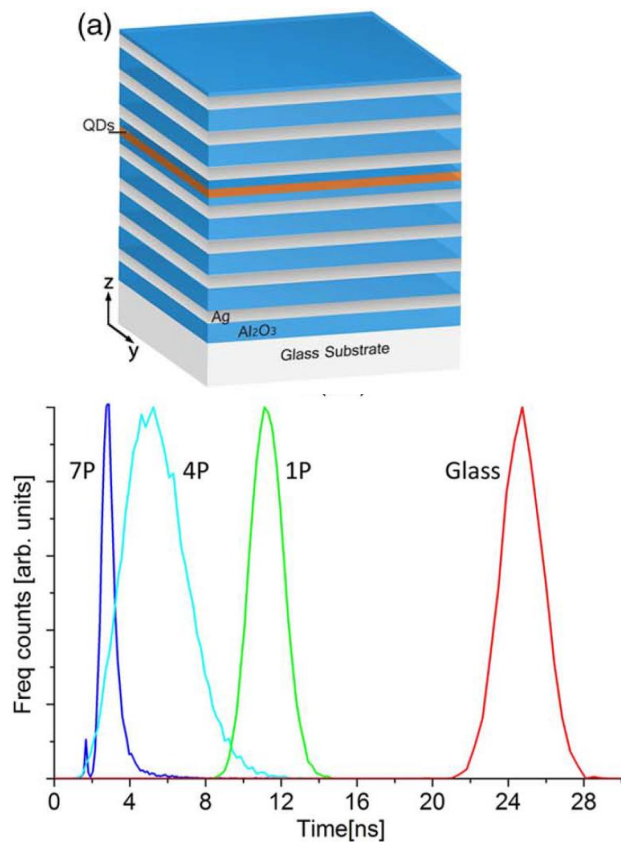


Liu, Z. et al., Science, vol. 315, 1686. (2007)

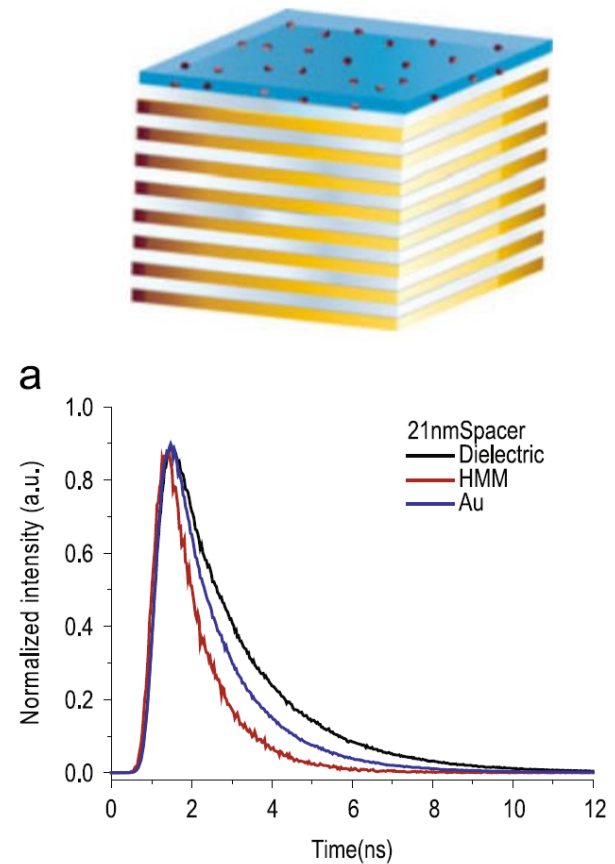
B. D. F. Casse et al., Appl. Phys. Lett., vol. 96, 023114 (2010)

Extremely high PDOS

Nonresonant phenomena → Broadband extremely high PDOS
Spontaneous emission engineering possible

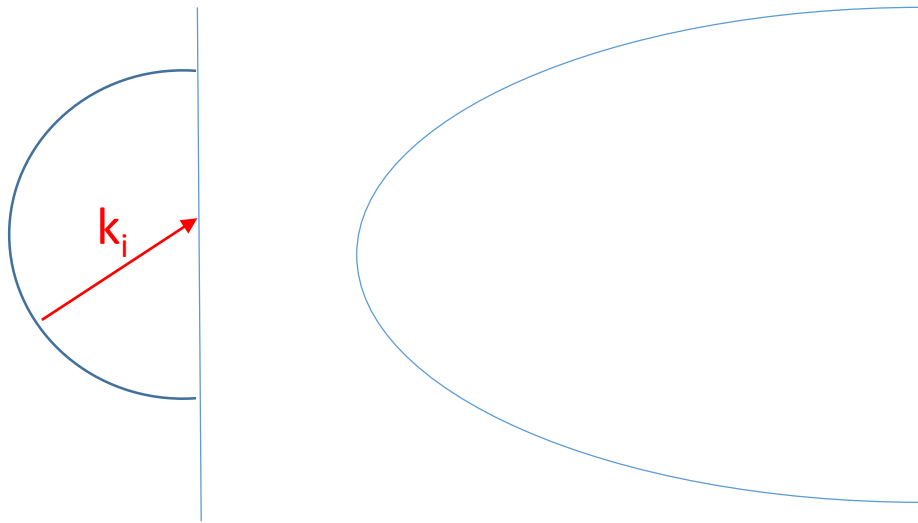
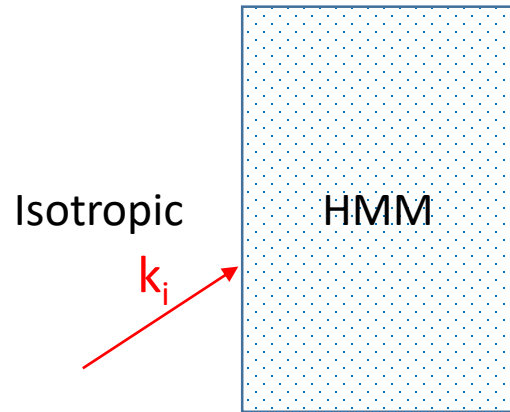


Galfsky, T. et al., *Optica*, vol. 2, 62-65. (2015)

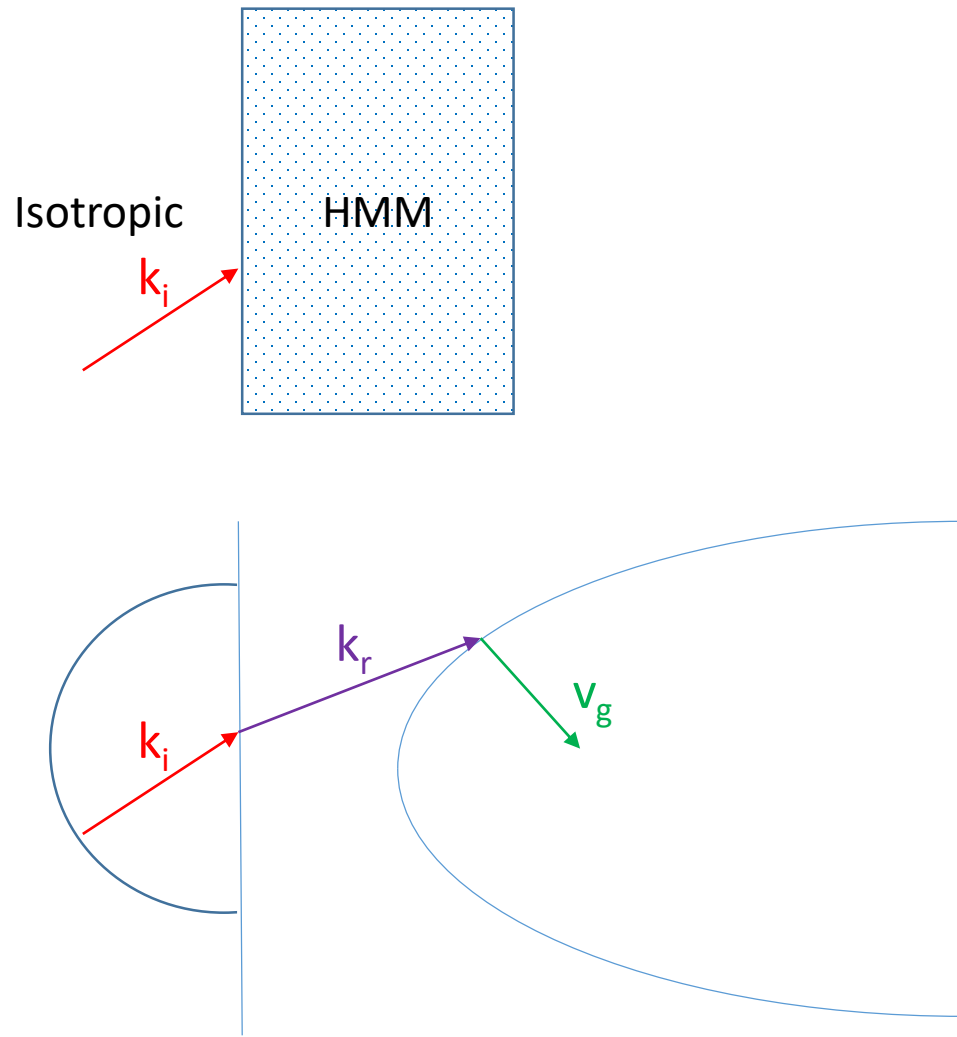


Z. Jacob et al, *Applied Physics B*, vol. 100, 215. (2010)

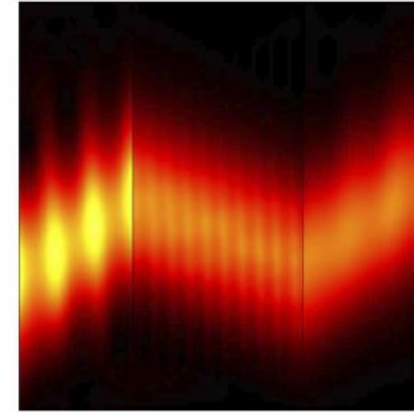
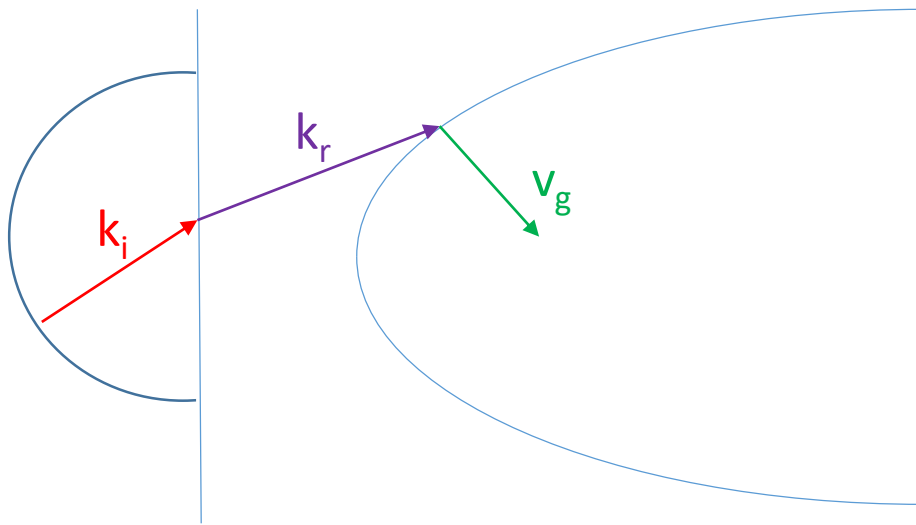
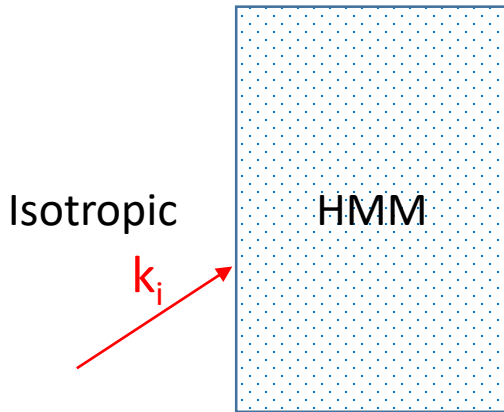
Negative refraction



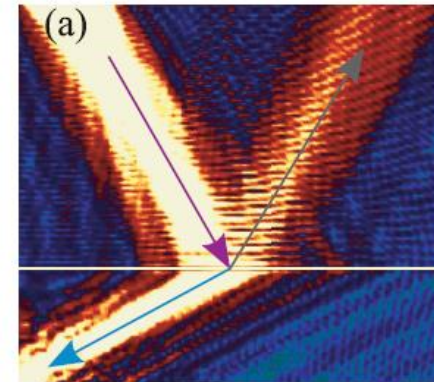
Negative refraction



Negative refraction



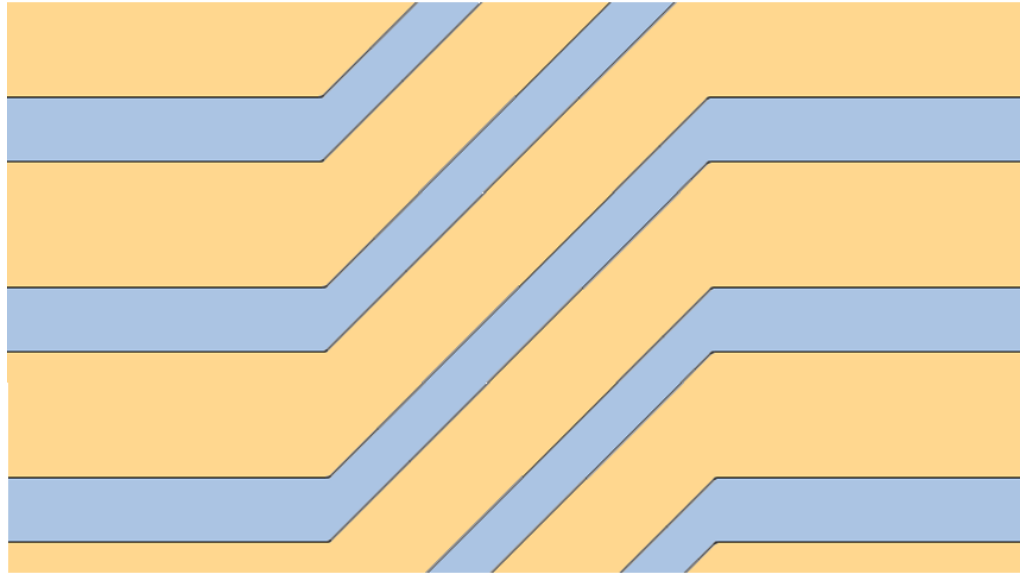
Y. Liu et al, Optics Express, vol. 16, 15439. (2008)



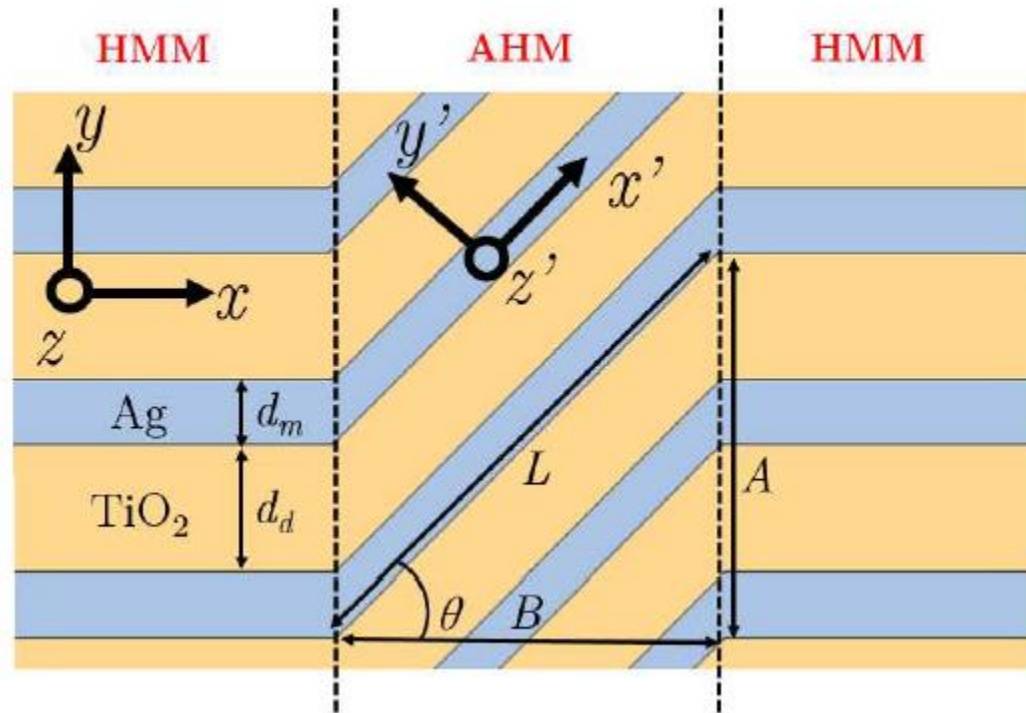
A. Orlov et al, Physical Review B, vol. 84, 045424 (2011)

- Introduction to hyperbolic metamaterials (HMMs)
- Some properties
- **Hyperbolic cavities with Fano resonances**

Reflection and transmission in slanted cavities



Reflection and transmission in slanted cavities



Right and left: simple multilayer HMM

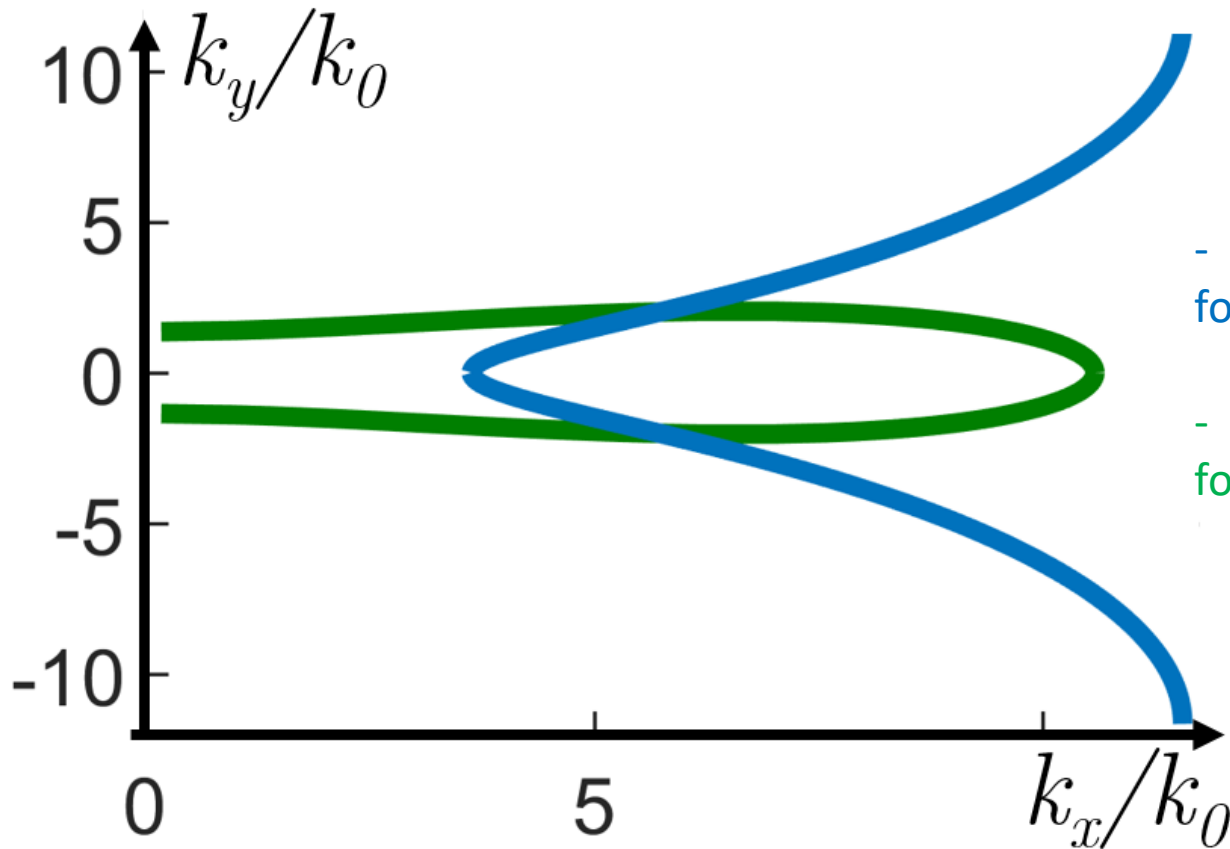
Centre: « asymmetric hyperbolic metamaterial » (tilted optical axis)

Exact solution (without losses in metal)

$$\cos(k_y D) = \frac{(\kappa_d \varepsilon_m + \kappa_m \varepsilon_d)^2}{4\kappa_d \kappa_m \varepsilon_d \varepsilon_m} \cosh(\kappa_d d_d + \kappa_m d_m) - \frac{(\kappa_d \varepsilon_m - \kappa_m \varepsilon_d)^2}{4\kappa_d \kappa_m \varepsilon_d \varepsilon_m} \cosh(\kappa_d d_d - \kappa_m d_m)$$

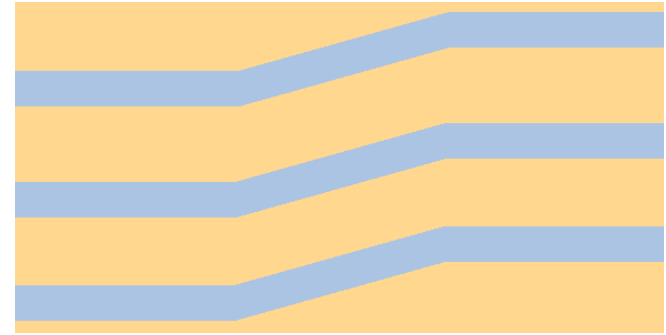
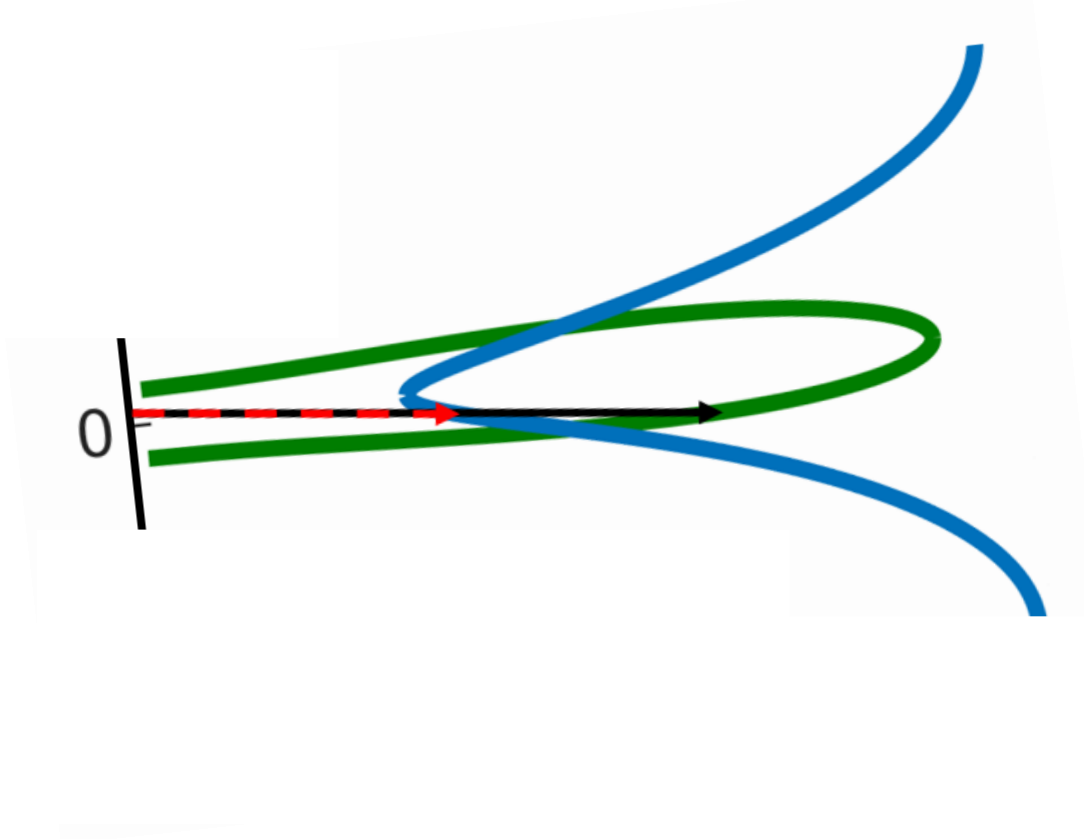
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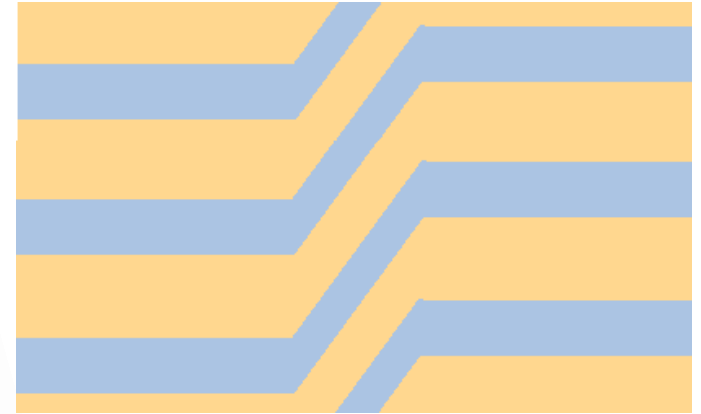
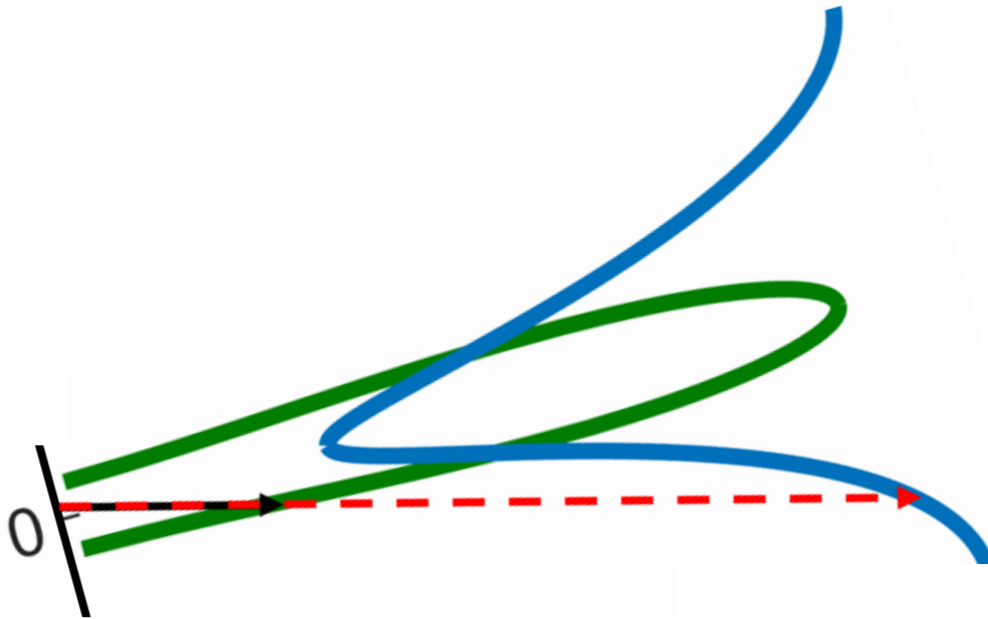
- Periodic isofrequency curve for the propagative mode
- Close isofrequency curve for the evanescent mode

Transverse momentum conservation ($k_y = 0$)



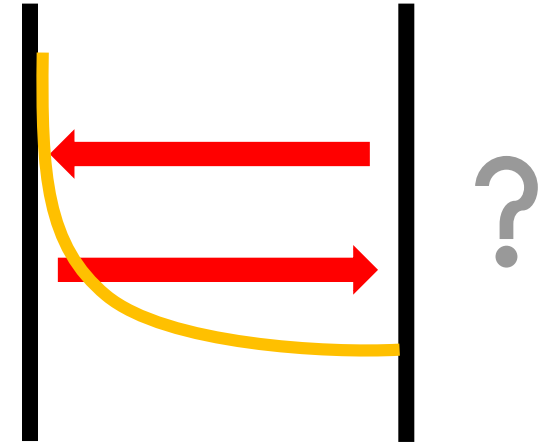
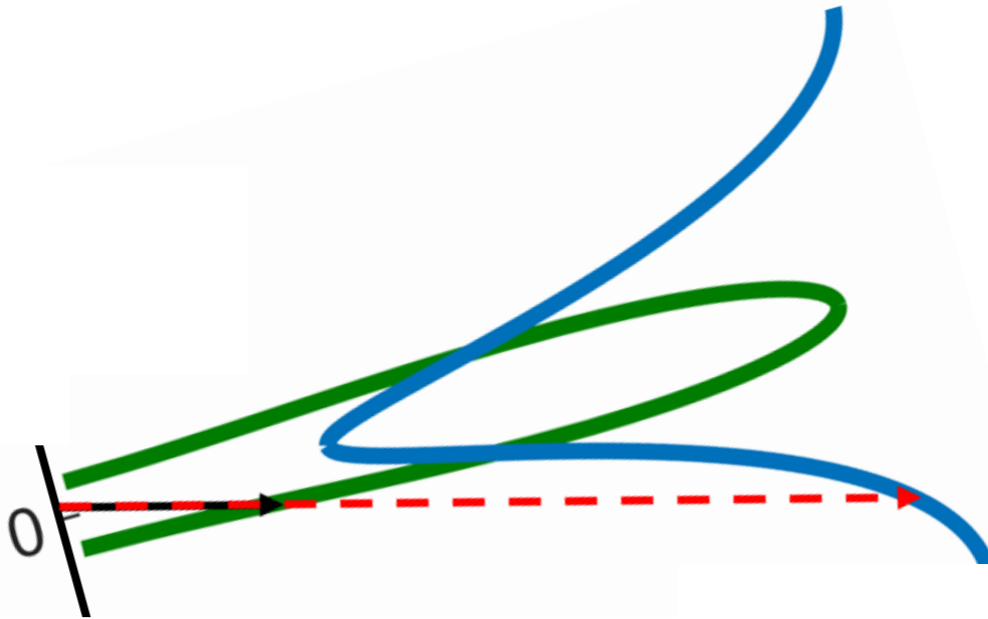
Always a propagative and evanescent mode excited !

Transverse momentum conservation



Always a propagative and evanescent mode excited !

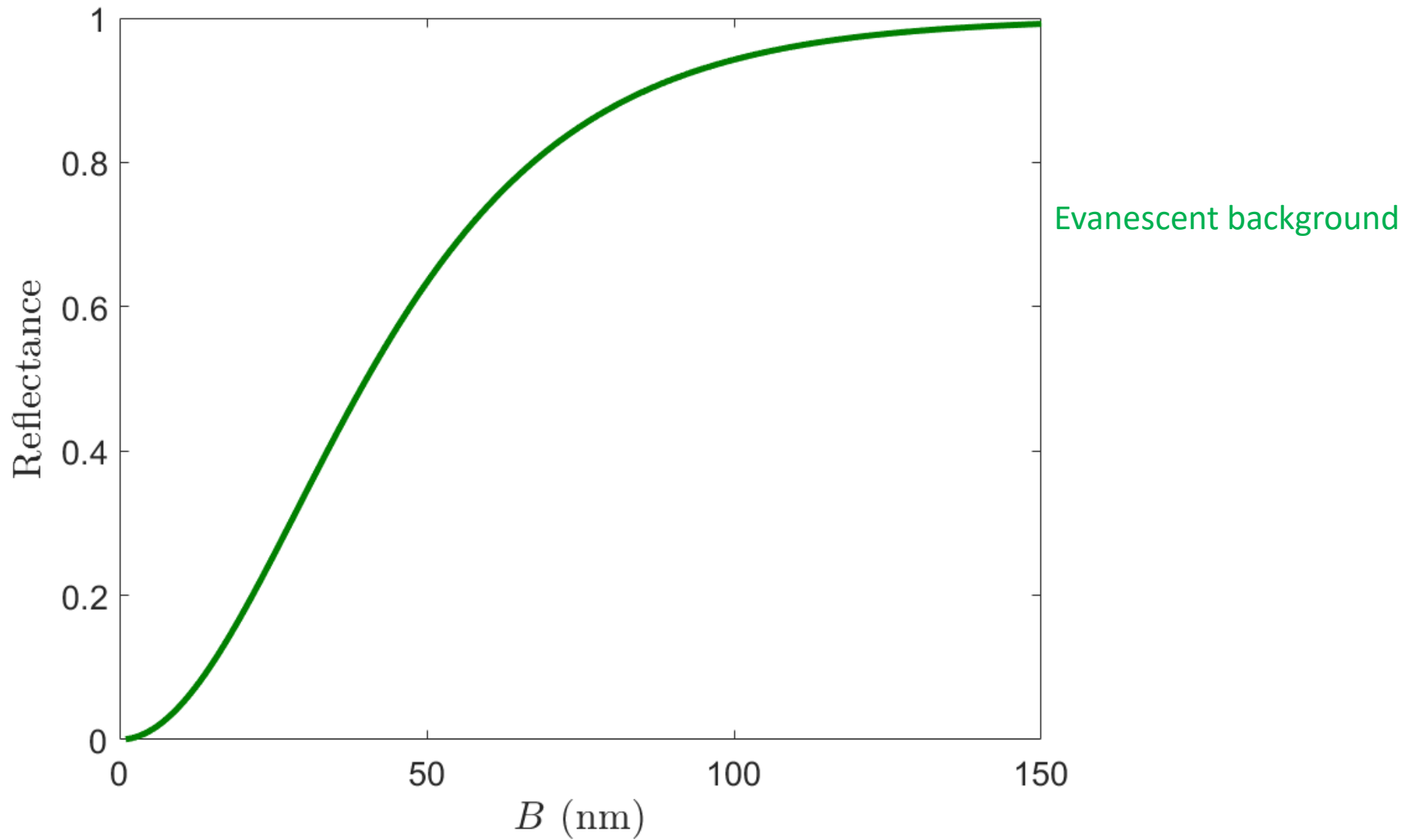
Transverse momentum conservation



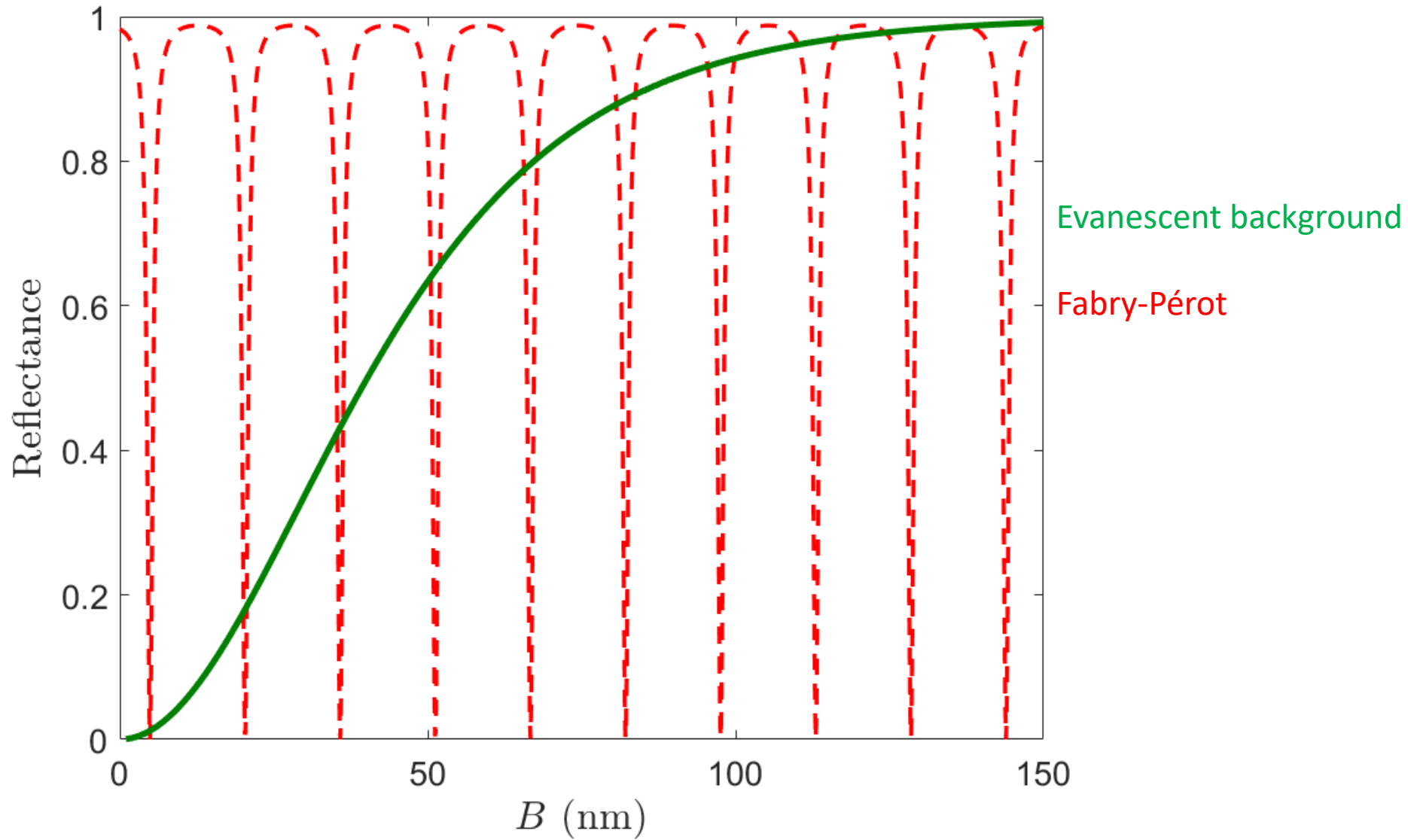
Always a propagative and evanescent mode excited !

→ Interference at the output

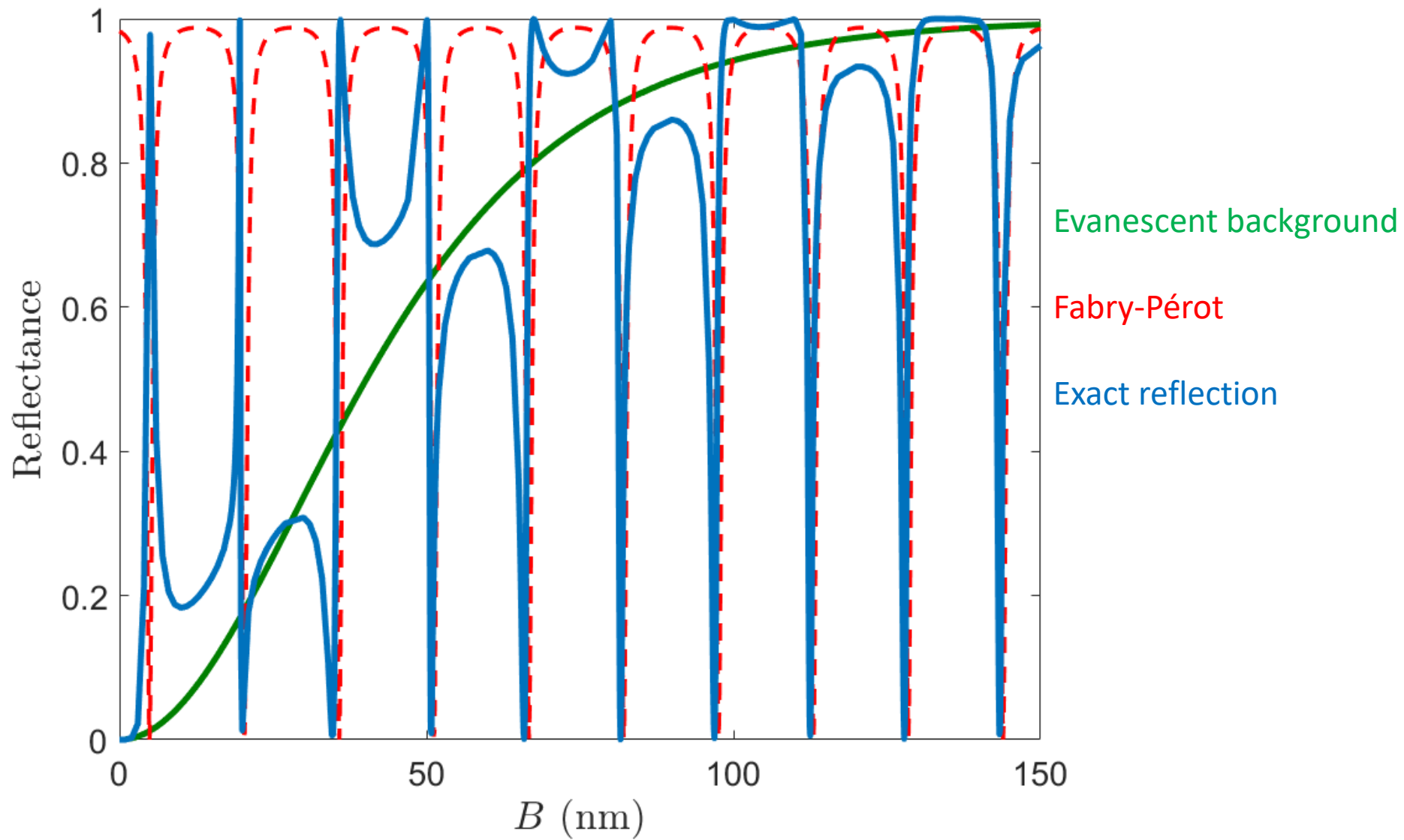
Fano resonances ($\Theta = 45^\circ$)



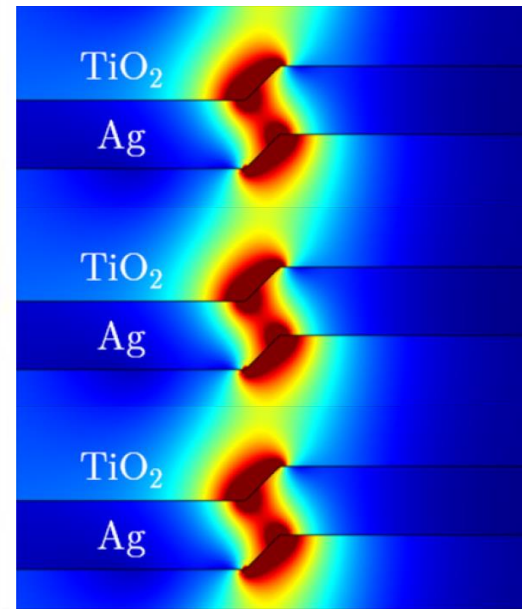
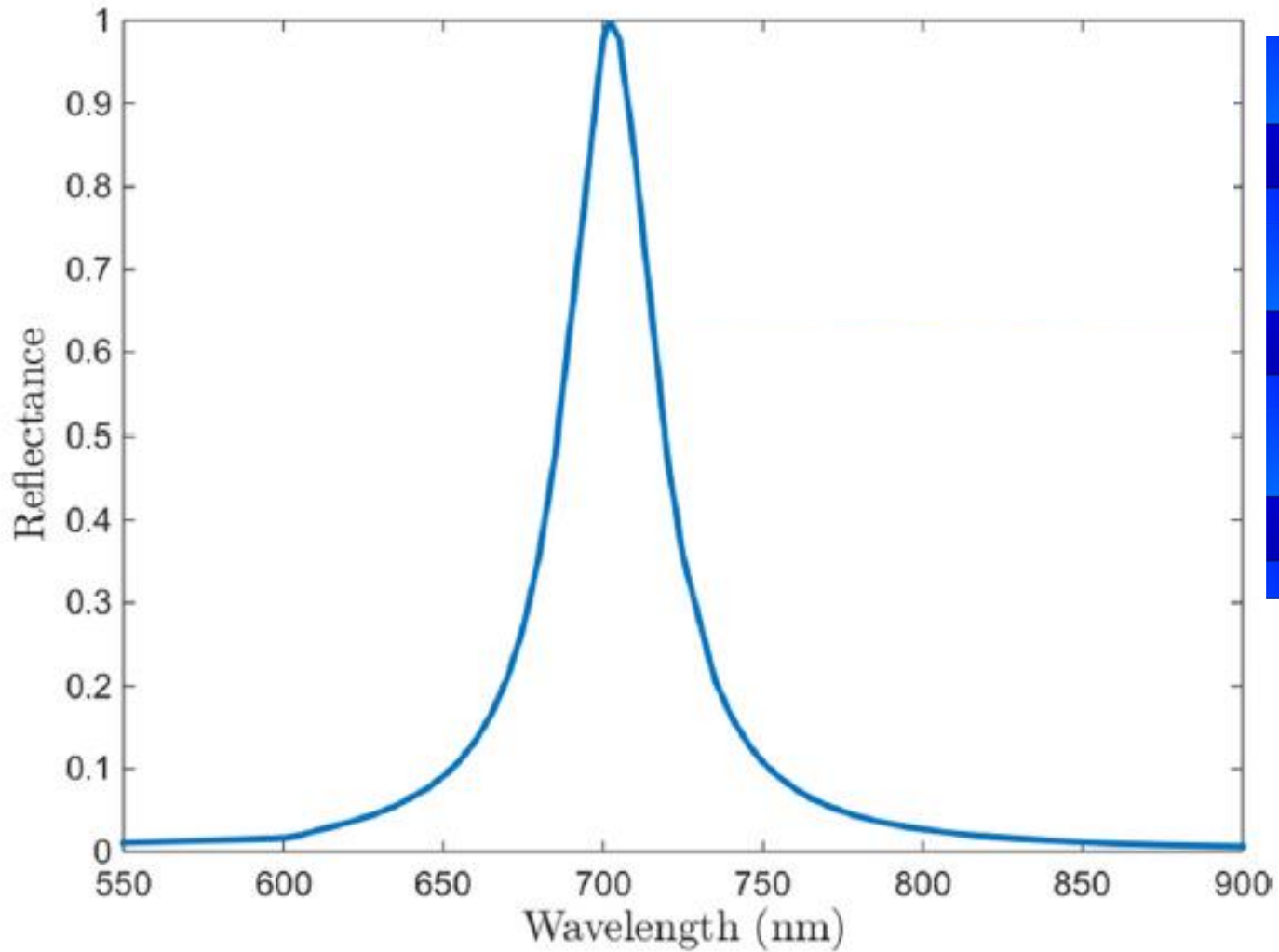
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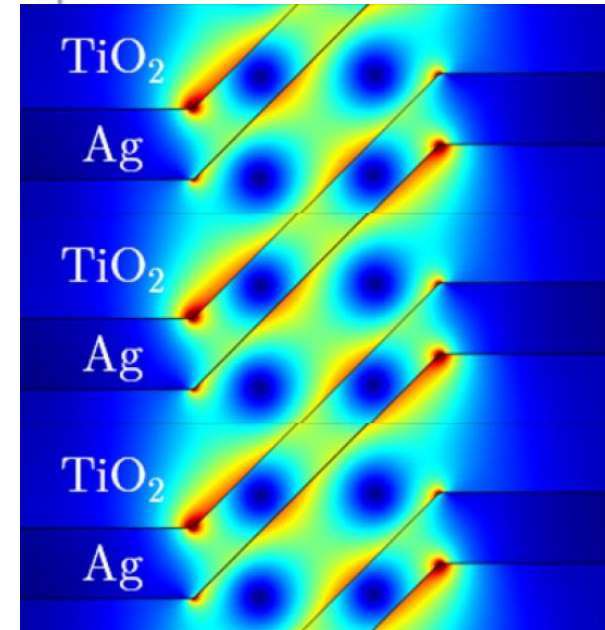
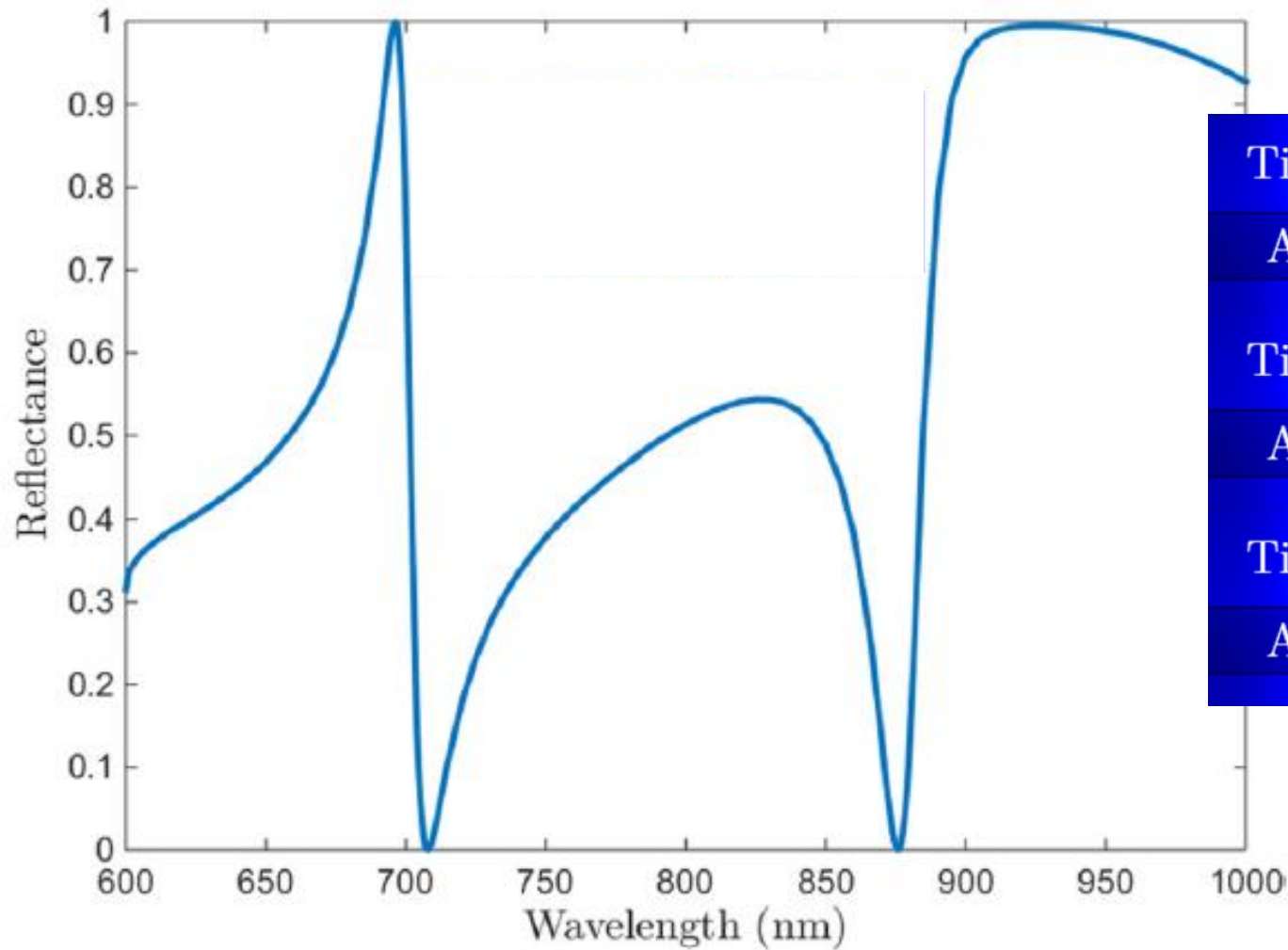
Fano resonances ($\Theta = 45^\circ$)



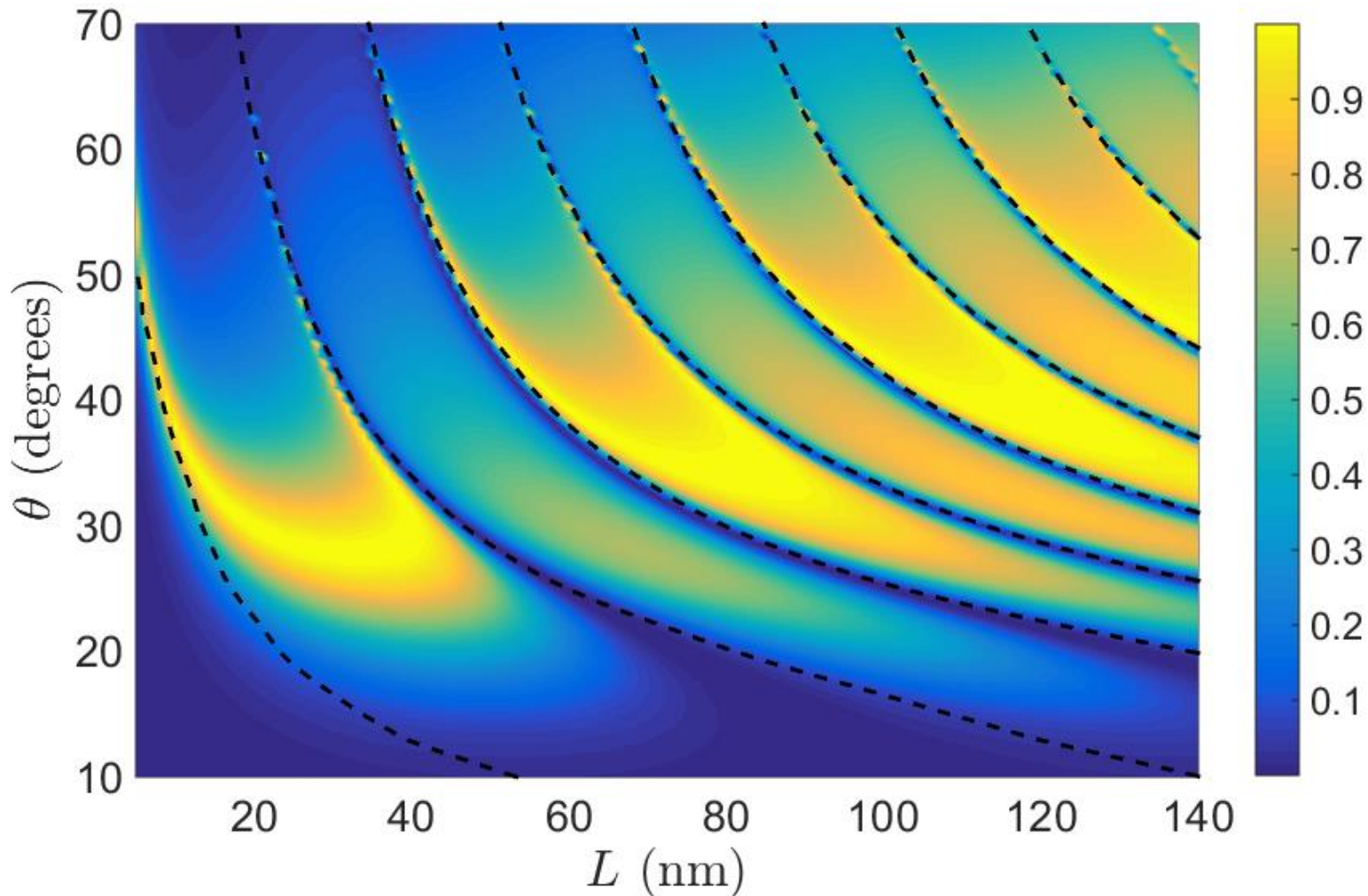
Spectrum for $B = 5$ nm



Spectrum for $B = 35$ nm



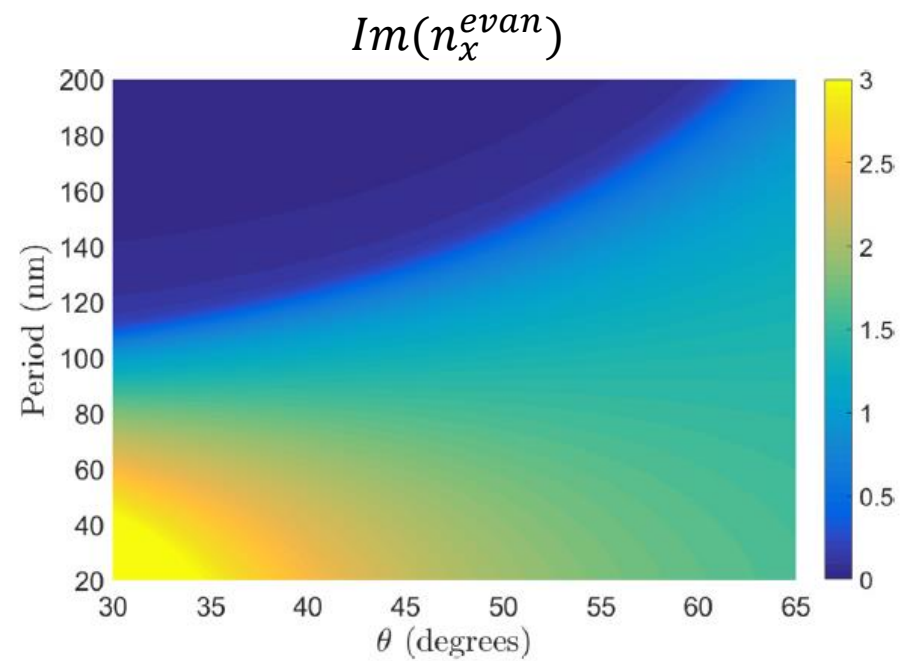
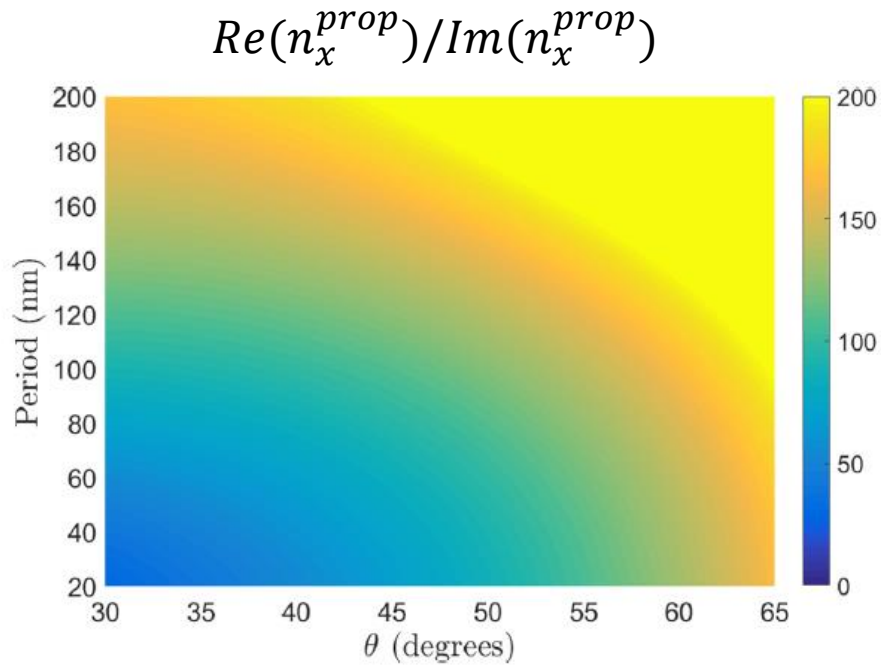
Reflection map (without loss)



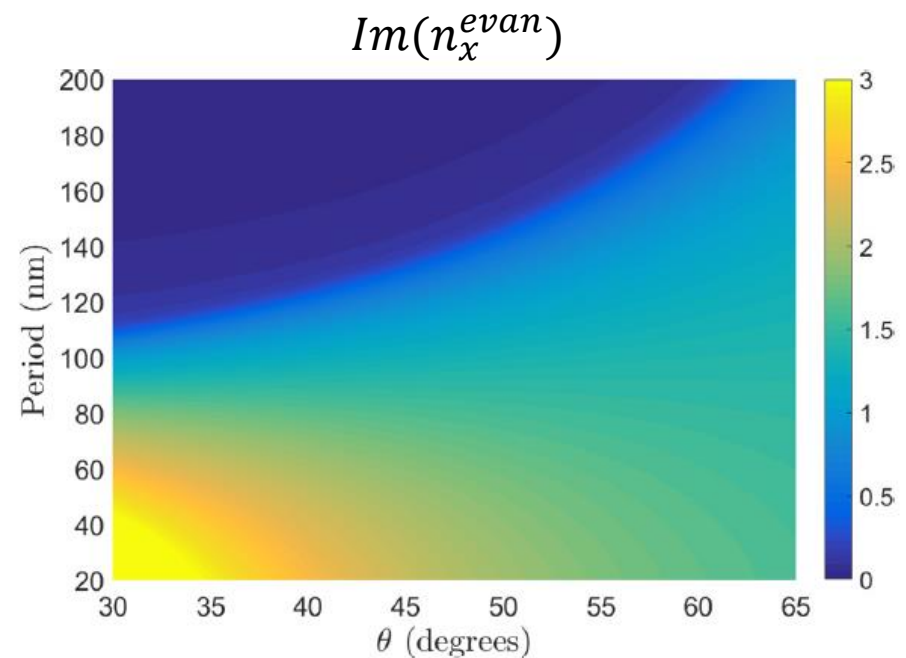
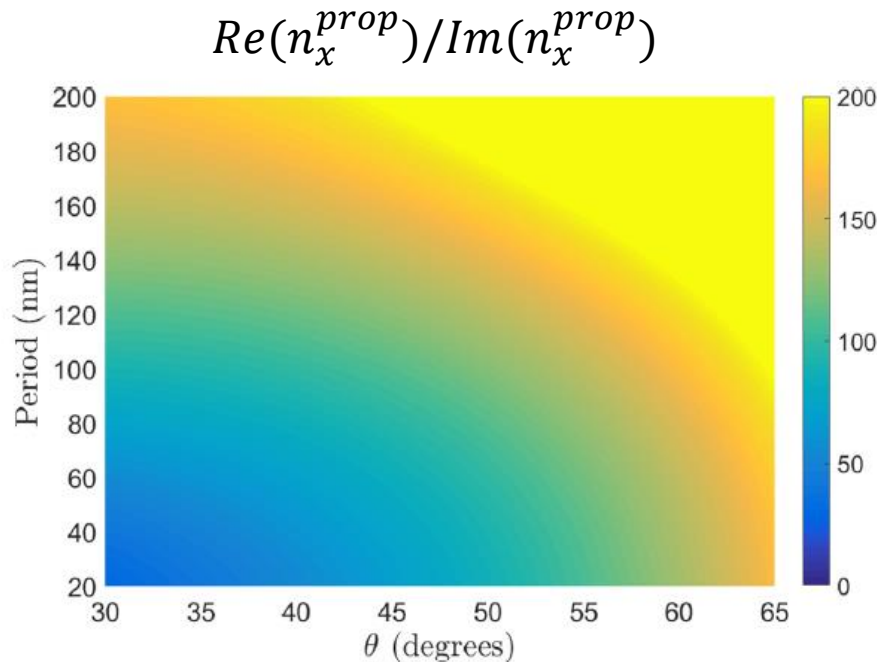
--- $2k_x(\theta)B + 2\varphi(\theta) = 2\pi m$ Phase matching

F. Vaianella and B. Maes, Physical Review B, vol. 94, pp 125442. (2016)

Lossy metal : condition for Fano resonances

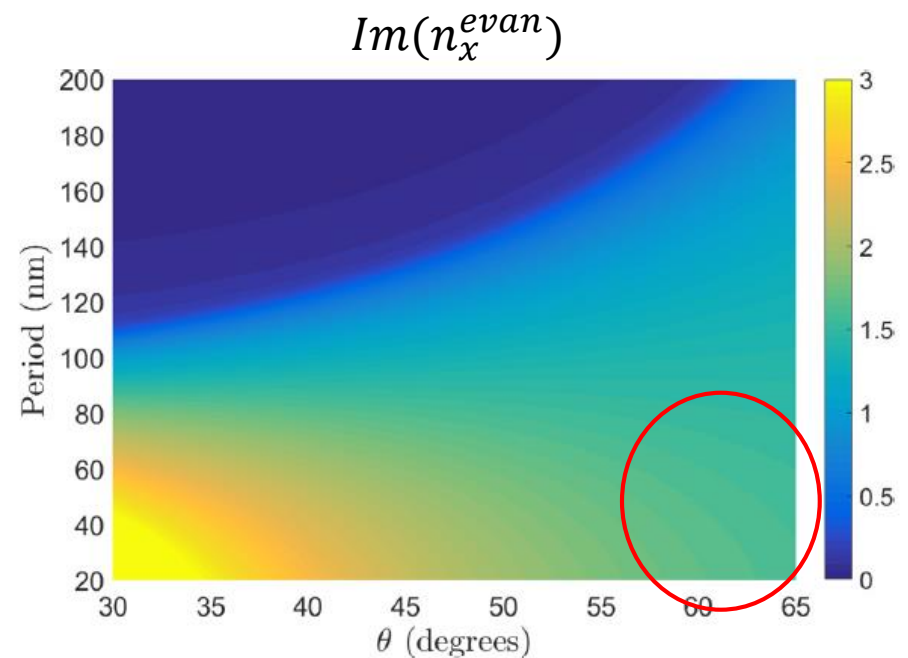
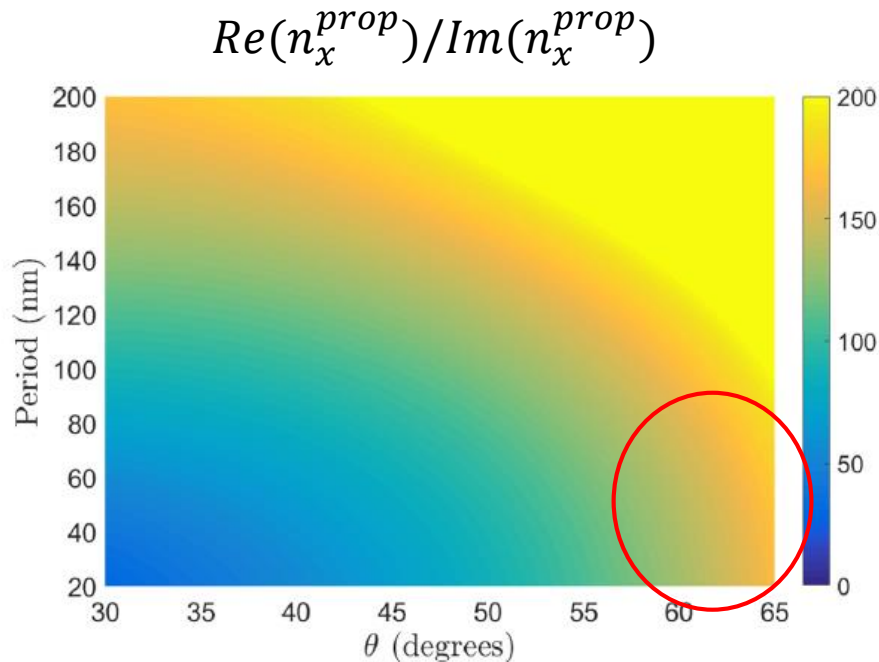


Lossy metal : condition for Fano resonances



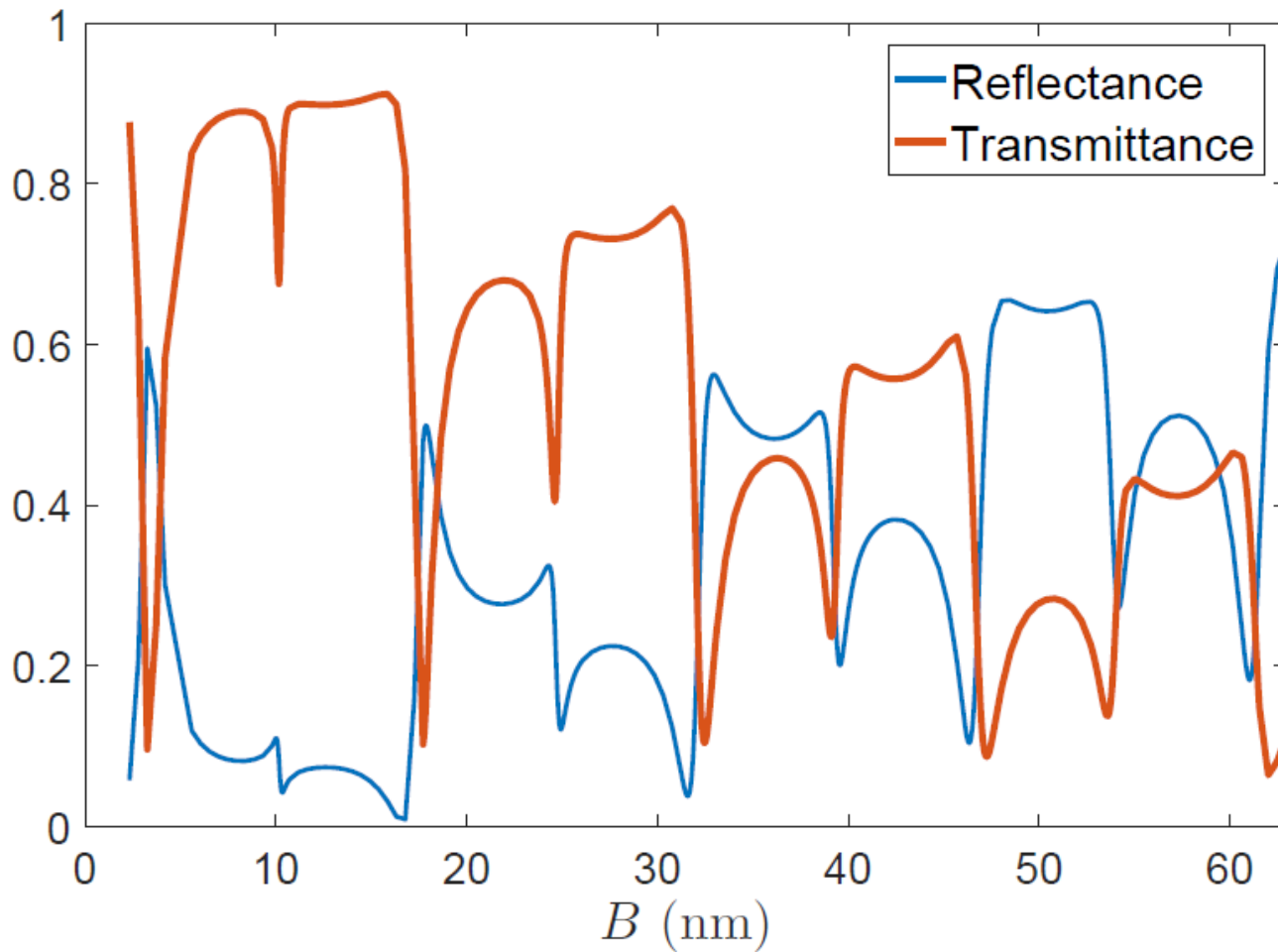
- Propagating mode should have large real part and small imaginary part of refractive effective index
- Evanescent mode should have imaginary part not too high (background would disappear) and not too low (background not efficient)

Lossy metal : conditions for Fano resonances



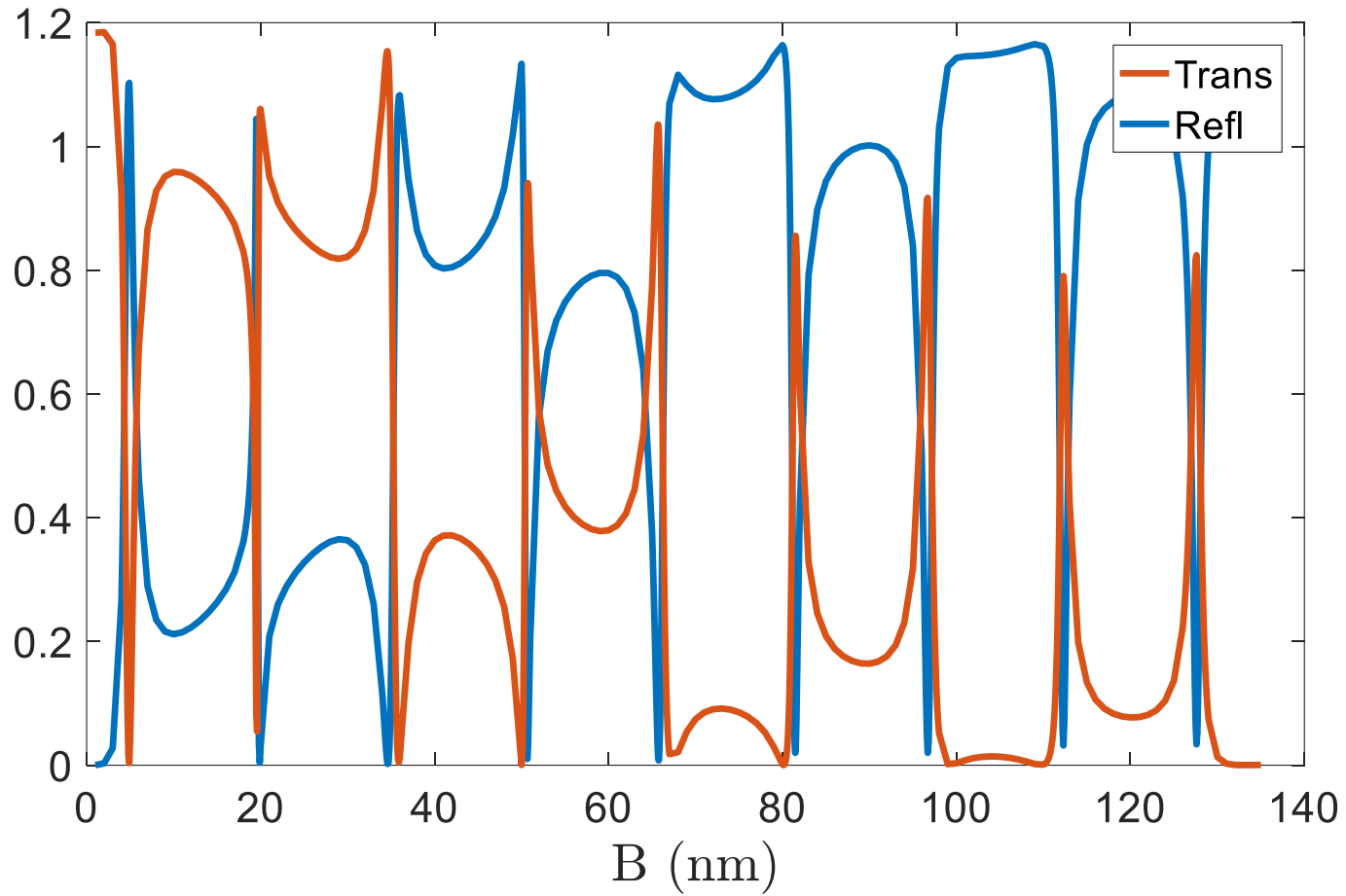
- Propagating mode should have large real part and small imaginary part of refractive effective index
- Evanescent mode should have imaginary part not too high (background would disappear) and not too low (background not efficient)

Scattering with losses for $\Theta = 65^\circ$

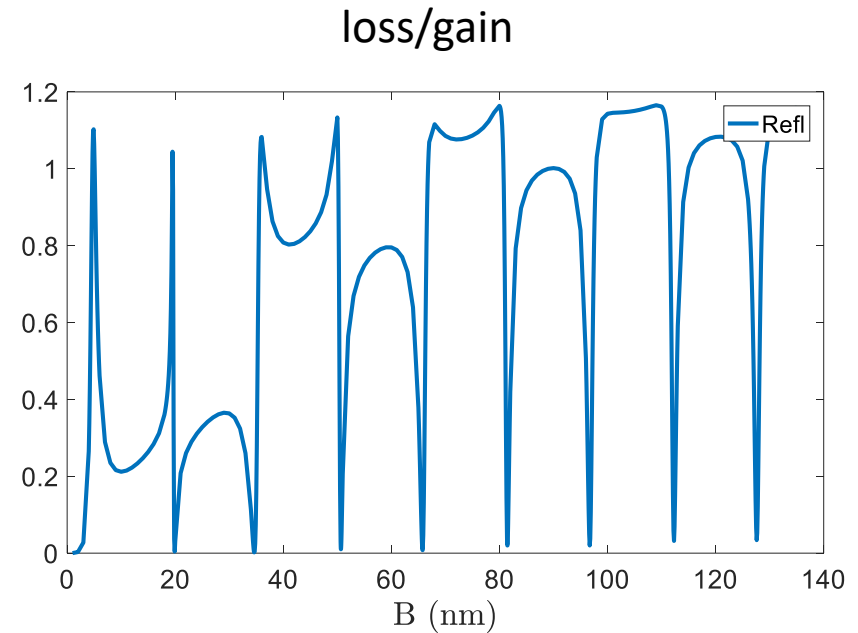
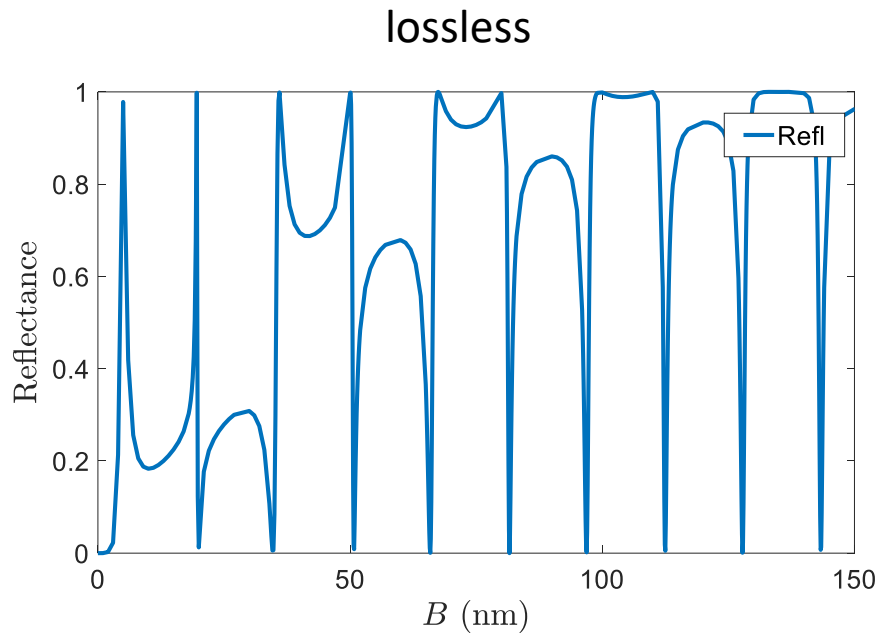


Fano resonances still present but more or less damped

Introduction of gain in the dielectric : $\text{Im}(n_{\text{TiO}_2}) = -0,07$



Comparison lossless – gain/loss structures



Introduction of gain allows 100% transmittance-reflectance Fano resonances
Actually difficult to introduce gain in TiO₂
Would be easier to work with semiconductors in infrared regime

Conclusions

- Hyperbolic metamaterials are periodic plasmonic structures with positive component of dielectric tensor in one direction and negative in another
- Fano resonances in ultra compact cavities for great control of the reflection and transmission of light
- Effective medium approximation inaccurate for this work. Predicts the excitation of one single mode, no Fano resonances possible
- Other topics: Heat transfer, active HMM, tunable HMM with graphene, homogenization theory, ...

Thank you for your attention

This work is financially supported by the F.R.I.A.-F.N.R.S.